

Proceedings of the First Workshop Organized by the IAFSS Working Group on Measurement and Computation of Fire Phenomena (MaCFP)

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Abstract

This paper provides a report of the discussions held at the first workshop on Measurement and Computation of Fire Phenomena (MaCFP) on June 10-11 2017. The first MaCFP workshop was both a technical meeting for the gas phase subgroup and a planning meeting for the condensed phase subgroup. The gas phase subgroup reported on a first suite of experimental-computational comparisons corresponding to an initial list of target experiments. The initial list of target experiments identifies a series of benchmark configurations with databases deemed suitable for validation of fire models based on a Computational Fluid Dynamics approach. The simulations presented at the first MaCFP workshop feature fine grid resolution at the millimeter- or centimeter-scale: these simulations allow an evaluation of the performance of fire models under high-resolution conditions in which the impact of numerical errors is reduced and many of the discrepancies between experimental data and computational results may be attributed to modeling errors. The experimental-computational comparisons are archived on the MaCFP repository [1]. Furthermore, the condensed phase subgroup presented a review of the main issues associated with measurements and modeling of pyrolysis phenomena. Overall, the first workshop provided an illustration of the potential of MaCFP in providing a response to the general need for greater levels of integration and coordination in fire research, and specifically to the particular needs of model validation.

Keywords: Buoyant plumes, Pool fires, Wall fires, Flame extinction, Fire modeling, Pyrolysis

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1. Introduction

A new initiative, endorsed and supported by the International Association for Fire Safety Science (IAFSS) [2], has been launched: the “*IAFSS Working Group on Measurement and Computation of Fire Phenomena*” (or the MaCFP Working Group) [3]. The general objective of the MaCFP Working Group is to establish a structured effort in the fire research community in order to make significant and systematic progress in fire modeling through a fundamental understanding of fire phenomena. The technical objectives are to develop the scientific foundations for the application of fire models to current or new challenging areas, for instance, flame spread, fire suppression, smoke toxicity. This is to be achieved as a joint effort between experimentalists and modelers on the general topic of the experimental validation of fire models based on a Computational Fluid Dynamics (CFD) approach. The MaCFP Working Group is intended as an open, community-wide, international collaboration between fire scientists. It is also intended to become a regular series of workshops, with workshops held every two or three years. The first workshop organized by the MaCFP Working Group was held on June 10-11 2017 as a pre-event to the 12th IAFSS Symposium in Lund, Sweden [4]. This paper presents a summary of the discussions and outcomes of the first MaCFP workshop.

The content and format of the first MaCFP workshop had been previously decided during a planning meeting in 2015. The planning meeting had produced a list of target experiments with databases deemed suitable for validation of CFD-based fire models. The intent was to make sure that the first workshop would go beyond the level of general discussions and would include presentations of a first suite of experimental-computational comparisons corresponding to an initial list of relevant experiments. The list of target experiments and a call for participation in the first workshop were broadly advertised to the international fire research community through letters to the editors of *Fire Safety Journal* [5] and *Fire Technology* [6] as well as through emails to the IAFSS membership.

While early discussions of the MaCFP Working Group had focused on gas phase phenomena (primarily flow and combustion phenomena), discussions were started in 2016 to expand the scope of MaCFP to include a subgroup dedicated to the modeling of pyrolysis phenomena. This led to a re-structuring of MaCFP into two subgroups: the (original) “*gas phase subgroup*” and the (new) “*condensed phase subgroup*”. Thus, in addition to being a first technical meeting for the gas phase

31 subgroup, the June-2017 MaCFP workshop also served as a planning meeting for the condensed
32 phase subgroup. The planning meeting portion of the workshop included a review of the main
33 issues associated with pyrolysis measurements and modeling for fire applications and a discussion
34 of future priorities for the condensed phase subgroup.

35 The technical meeting portion of the workshop provided a first demonstration of current activ-
36 ities of the MaCFP Working Group as well as an illustration of their potential impact. The initial
37 list of target experiments identified by the gas phase subgroup corresponds to basic configurations
38 (also called building blocks) featuring carefully-controlled conditions and quality instrumentation
39 and diagnostics. They also correspond to experiments with open, easily-accessible databases. In
40 what is considered as a first intermediate step, the list has a limited scope and only includes simple
41 turbulent buoyant plumes and simple flames (in most cases, the flames are non-sooting or only
42 weakly-sooting), supplied with gaseous or liquid fuel, and featuring open burn conditions; the case
43 of strongly sooting and smoking flames, fueled by solid flammable materials, and featuring com-
44 partment effects is outside the scope of the first MaCFP workshop and will be considered in future
45 editions.

46 The initial list of MaCFP target experiments includes five categories:

- 47 • (Case 1) Turbulent buoyant plumes: this category corresponds to open plumes and is repre-
48 sented by a helium plume experiment conducted at Sandia National Laboratories (Sandia) [7].
- 49 • (Case 2) Turbulent pool fires with gaseous fuel: this category corresponds to open flames with
50 a prescribed fuel flow rate and is represented by a series of natural gas flame experiments
51 conducted at the National Institute of Standards and Technology (NIST) [8] (and referred to in
52 the following as the NIST McCaffrey natural gas flame experiment) and by a series of methane
53 and hydrogen fire experiments conducted at Sandia National Laboratories (Sandia) [9, 10].
- 54 • (Case 3) Turbulent pool fires with liquid fuel: this category corresponds to open flames with a
55 thermal-feedback-driven fuel flow rate and is represented by a methanol pool fire experiment
56 conducted at the University of Waterloo (UW) [11, 12].
- 57 • (Case 4) Turbulent wall fires: this category corresponds to boundary layer flames with a
58 prescribed fuel flow rate and is represented by a series of vertical wall flame experiments,
59 fueled by methane, ethane, ethylene or propylene, and conducted at FM Global [13, 14].
- 60 • (Case 5) Flame extinction: this category corresponds to flames driven to extinction conditions

61 and is represented by a series of methane and propane line flame experiments conducted at
62 the University of Maryland (UMD) [15–17].

63 Note that the experimental databases corresponding to Cases 1-5 are hosted on the MaCFP
64 repository [1] with open access so that the data are available to the fire research community as
65 reference data for future experimental and/or computational studies.

66 Seven groups submitted computational results for comparisons with experimental data and for
67 discussions at the first MaCFP workshop: FM Global (USA); Ghent University (UGent, Belgium);
68 the Institut de Radioprotection et de Sûreté Nucléaire (IRSN, France); the National Institute of
69 Standards and Technology (NIST, USA) teamed up with the VTT Technical Research Centre
70 of Finland (VTT, Finland); Sandia National Laboratories (SNL, USA); University of Cantabria
71 (UCantabria, Spain); and University of Maryland (UMD, USA). These groups used one of the
72 following four CFD solvers:

- 73 • FDS (Fire Dynamics Simulator) developed by NIST in collaboration with VTT [18];
- 74 • FireFOAM based on OpenFOAM [19] and developed by FM Global [20];
- 75 • ISIS developed by IRSN [21];
- 76 • SIERRA/Fuego developed by SNL [22].

77 These solvers are representative of current fire modeling capabilities available for research-level
78 and/or engineering-level projects.

79 The paper is structured as follows. Section 2 describes the main outcomes of the technical
80 meeting held by the gas phase subgroup of the MaCFP Working Group. Section 2.1 gives a
81 brief description of the different concepts used in quality control of CFD models and a review
82 of the computational challenges found in model validation (the focus of MaCFP). Sections 2.2-
83 2.6 present a summary of the experimental-computational comparisons performed for Cases 1-5,
84 respectively. Section 2.7 presents a conclusion and a description of future plans for the gas phase
85 subgroup. Section 3 presents a review of the discussions held during the planning meeting of the
86 condensed phase subgroup of the MaCFP Working Group. Section 3.1 gives a brief description
87 of the objectives of the subgroup. Section 3.2 presents a summary of the invited presentations
88 and follow-up discussion that took place at the workshop. Section 3.3 presents a conclusion and a
89 description of future plans for the condensed phase subgroup.

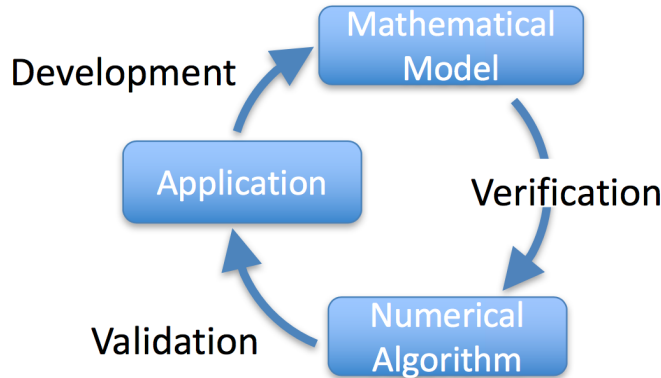


Figure 1: Schematic representation of the verification and validation process used to evaluate the accuracy of a computational model. Adapted from [26].

90 2. Gas Phase Subgroup

91 2.1. Different Aspects of Quality Control in CFD: Verification and Validation, Grid Resolution, 92 Physical Modeling

93 In this section, we first briefly put the current MaCFP effort in the general context of CFD
94 verification and validation (section 2.1.1). We then proceed to review the computational challenges
95 associated with simulations of the target experiments selected for the first MaCFP workshop. The
96 challenges include the design of the computational grid (section 2.1.2) and the uncertainties associ-
97 ated with model descriptions of turbulence, combustion and radiation phenomena (section 2.1.3).

98 2.1.1. Code Verification and Model Validation

99 The verification and validation (or V&V) process is the primary quality control method used
100 to establish the degree of confidence in a computational model for a specific application [23–29].
101 Code *verification* is the process of determining whether the model has been correctly implemented
102 on the computer. In other words, verification “checks the math”. Model *validation* is the process
103 of determining whether the model correctly represents the physical phenomena of interest. In other
104 words, validation “checks the physics”. The process of developing a complex fire model is a cycle,
105 as depicted in Fig. 1, whereby: (1) a mathematical model is proposed; (2) the model is verified;
106 (3) the model is validated; and finally, if modifications to the model are required to achieve more
107 accurate results for the intended application, then the process is repeated.

108 The prevailing technique for verification consists in comparing results of the computational
109 model with analytical solutions (or manufactured solutions [26]) obtained for the same system of
110 governing equations. This is usually accomplished through the construction of a library of *unit*

111 *test* problems. These problems are selected to exercise certain parts of the code (for instance the
112 flow solver — with or without convection, with or without diffusion, — the combustion solver, the
113 radiation solver, *etc*), considered sequentially and in isolation. The library is constructed with the
114 objective to attain as much “code coverage” with unit test problems as possible. Note that MaCFP
115 is not concerned with verification and assumes that the CFD solvers selected for MaCFP activities
116 have been and are continuously verified. MaCFP is focused on validation.

117 The prevailing technique for validation consists in comparing experimental data with results of
118 the computational model obtained in the same configuration. The target experiments considered in
119 MaCFP correspond to “open” validation tests in which the modelers have unlimited access to the
120 details of the setup and to the experimental data prior to running the model. In most simulations
121 presented below, the CFD solvers are used in their baseline configuration (presented briefly in
122 section 2.1.3), *i.e.*, without resorting to any model modification or calibration, and the simulations
123 can be therefore interpreted as true *validation* tests. In the case of the UMD turbulent line flame
124 experiments, the CFD solvers were used with some advanced features to describe flame extinction
125 that may or may not be part of the baseline configurations. In addition, some of these advanced
126 features were originally tuned against data obtained from the same UMD experiments. In that
127 case, the simulations should be interpreted as *calibration* tests rather than validation tests.

128 In this first edition of the MaCFP workshop series, no effort was made to impose particular
129 metrics in the comparison between experimental data and simulation results. Comparisons gener-
130 ally take the form of plotting measured and simulated spatial profiles of mean or root-mean-square
131 (*rms*) quantities (mean quantities refer to time-averaged quantities and *rms* quantities designate
132 the square root of the mean squared deviation of a quantity from its mean). Also no effort was
133 made to require systematic estimates of experimental or numerical uncertainties.

134 2.1.2. Computational Grid Design

135 One of the main challenges found in the application of CFD tools to the simulation of complex
136 flow problems is the design of the computational grid. In large eddy simulations (LES), the design of
137 the computational grid comes from an analysis of the characteristic length scales of the problem and
138 a requirement that large-scale features that (presumably) control the flow dynamics be captured
139 by the grid. The implicit assumption is that small-scale features are dynamically controlled by
140 the resolved scales and can be represented through subgrid-scale (SGS) models. Relevant large-
141 scale features that are considered dynamically-controlling and are therefore grid-resolved in LES

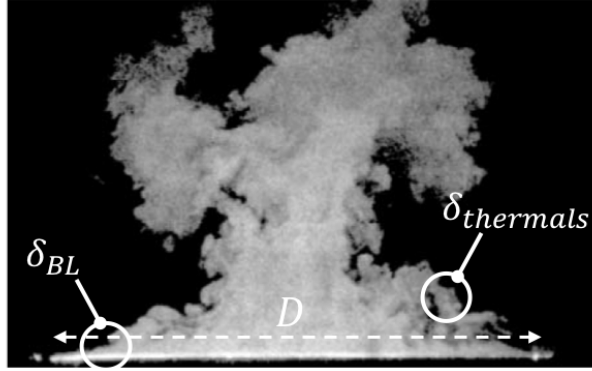


Figure 2: Illustration of the multi-scale nature of pool fire configurations using a flow visualization of the Sandia helium plume experiment. The configuration features: large-scale structures in the center of the plume with size on the order of the plume diameter D ; thin boundary layers near the edges of the plume with size δ_{BL} ; and small structures created by Rayleigh-Taylor instabilities with size $\delta_{thermals}$. Adapted from [7].

142 of turbulent diffusion flames include: the large flow structures responsible for the production of
 143 turbulent kinetic energy (their size is estimated by the integral length scale of the turbulent flow);
 144 the large wrinkles on the flame surface responsible for enhanced fuel-air mixing and heat release; and
 145 the large soot-containing structures inside and outside the flame zone responsible for flame emission
 146 and smoke absorption properties. Small-scale features that are considered dynamically-controlled
 147 and remain therefore grid-unresolved in LES include: the small flow structures responsible for the
 148 dissipation of turbulent kinetic energy (their size is estimated by the Kolmogorov length scale); the
 149 thin reactive layers that make up the micro-structure of a turbulent flame; and the thin elongated
 150 soot layers that make up the micro-structure of the flame radiation field.

151 While the discussion above provides a valuable framework, it is important to emphasize that
 152 the separation between large scales that are dynamically-controlling and small scales that are
 153 dynamically-controlled is somewhat artificial and is not necessarily obvious. For instance, let us
 154 consider the case of a simple pool flame fueled by a liquid chemical supplied through a circular
 155 burner of diameter D . The pool fire literature suggests that D is the only relevant length scale of the
 156 problem: the mean flame vertical height, the mean flame horizontal thickness and the characteristic
 157 size of the large turbulent flow structures expected in the flame region are all proportional to D
 158 and can be captured in a LES simulation provided that the grid spacing is 10-20 times smaller than
 159 D .

160 However, this is not the whole story. In many cases, the pool flame features a strong buoyancy-
 161 driven instability (called the puffing instability) that results in large oscillations in the (horizontal)

162 entrained air flow and the (vertical) combustion products flow. Under strongly unstable conditions,
163 the instantaneous flame takes different shapes during the instability cycle, including the shape of a
164 somewhat unexpected thin (horizontal) boundary layer flame established close to the pool surface
165 and produced by large peak values of the air flow velocity. The presence of the intermittent bound-
166 ary layer flame is generally over-looked in pool fire studies but circumstantial evidence suggests
167 that it plays an important dynamical role in the instability cycle and consequently needs to be cor-
168 rectly captured by the computational grid. Because its thickness δ_{BL} is much smaller than the pool
169 diameter D , the simulation of the intermittent boundary layer flame brings stringent constraints
170 to the design of the computational grid.

171 A related topic is the possible development of Rayleigh-Taylor instabilities in regions of the flame
172 with unstable thermal stratification and the associated formation of small plume-like structures,
173 often called “thermals”. When present, these thermals are believed to be responsible for enhanced
174 fuel-air mixing and heat release. Because their characteristic size $\delta_{thermals}$ is much smaller than
175 the pool diameter D , the simulation of the thermals brings additional stringent constraints to the
176 design of the computational grid.

177 Figure 2 presents an illustration of the multi-scale nature of pool fire configurations (albeit in
178 the case of a chemically inert plume) and identifies the three dynamically-**important** length scales:
179 D , δ_{BL} and $\delta_{thermals}$. The dynamic effects occurring at these length scales can be captured in a LES
180 simulation provided that the grid spacing is 10 times smaller than D , δ_{BL} and $\delta_{thermals}$. **The choice**
181 **of a sufficiently fine computational grid that directly captures dynamical effects at all relevant**
182 **length scales may produce the best results but also corresponds to a high (or even prohibitive)**
183 **computational cost. Alternatively, the dynamic effects in pool fires can also be captured in a LES**
184 **simulation provided that the grid spacing is 10-20 times smaller than D and that subgrid-scale**
185 **models correctly represent the effects occurring at scales δ_{BL} and $\delta_{thermals}$.**

186 We now consider the implications of the previous discussion to the choice of grid resolution in
187 LES simulations of the target experiments selected for the first MaCFP workshop. In the list of
188 six target experiments (Cases 1-5 in section 1), **four (Cases 1-3)** correspond to pool-like configura-
189 tions with significant buoyancy-driven instability phenomena (*i.e.* strong puffing motions and the
190 formation of thermals): the Sandia helium plume; **the NIST McCaffrey natural gas flames**; the
191 Sandia methane and hydrogen **gas flames**; and the UW methanol pool fire. **The Sandia, NIST and**
192 **UW experiments feature pool diameters between 0.3 and 1 m, which based on the discussion above,**
193 **suggests that a centimeter-scale computational grid is required for simulations aimed at resolving D**

194 while relying on subgrid-scale models to represent the effects occurring at scales δ_{BL} and $\delta_{thermals}$.
195 In addition, the flow visualization techniques used in the Sandia and UW experiments reveal inter-
196 mittent boundary layers and thermals with length scales on the order of 1 cm, which suggests that
197 a millimeter-scale computational grid may be required for simulations aimed at resolving δ_{BL} and
198 $\delta_{thermals}$. Furthermore, the list of experiments selected for the first MaCFP workshop features two
199 additional configurations. One target experiment (Case 4) corresponds to a boundary layer flame
200 configuration: the FM Global vertical wall flame. The flow visualization techniques used in this
201 experiment reveal boundary layers and structures with length scales on the order of 1 cm, which
202 suggests using a millimeter-scale computational grid. The other target experiment (Case 5) corre-
203 sponds to a pool-like configuration without reported evidence of strong unstable motions: the UMD
204 methane and propane turbulent line flames. In the case of the UMD experiments, the only appar-
205 ent characteristic length scale is the burner width (5 cm), which suggests using a millimeter-scale
206 computational grid.

207 Note that the numerical submissions to the first MaCFP workshop generally correspond to grid
208 resolutions consistent with the estimates above. Submissions for the FM Global vertical wall flame
209 used similar levels of grid resolution (millimeter-scale); the design of the computational grid was
210 guided by the objective to correctly resolve the thickness of the boundary layer flame. Submissions
211 for the UMD turbulent line flames also used similar levels of grid resolution (millimeter-scale); the
212 design of the computational grid was guided by the objective to correctly resolve the thickness of
213 the line flame. In contrast, submissions for the pool-like configurations (the Sandia helium plume,
214 the NIST McCaffrey flames, the Sandia gas flames and the UW pool fire) varied more substantially
215 (from millimeter- to centimeter-scale; the pool-diameter-to-cell-size ratio varied between 20 and 120);
216 the design of the computational grid was guided by the objective to correctly resolve the diameter
217 of the pool and also (for simulations with the finest grids) by the objective to capture some of the
218 smaller-scale effects associated with δ_{BL} and $\delta_{thermals}$. In this first edition of the MaCFP workshop
219 series, no effort was made to require specific levels of resolution or systematic grid convergence
220 studies.

221 *2.1.3. Physical Modeling*

222 The numerical submissions to the first MaCFP workshop correspond to one of the following four
223 CFD solvers: FDS, FireFOAM, ISIS and SIERRA/Fuego. All four models are used in LES mode.
224 The solvers differ in details of the formulation of the governing equations, in the construction of the

225 computational grid, in the choice of algorithms used to discretize the governing equations and to
226 provide a numerical solution, and in their ability to perform parallel computing. These differences
227 are believed to be inconsequential in the present tests because the simulations correspond to simple
228 academic configurations. The solvers also differ in the formulation of the physical models used to
229 describe subgrid-scale turbulence, combustion and radiation. These differences are believed to be
230 significant and may be responsible for some or many of the reported discrepancies observed between
231 simulation results, as summarized in the following sections.

232 In their baseline configuration, the four CFD solvers use:

- 233 • For subgrid-scale turbulence: a classical gradient transport formulation with a SGS turbulent
234 viscosity. A closure expression for the SGS turbulent viscosity is provided by: a modified
235 Deardorff model [30] (FDS, see also Ref. [31]); a model using a (constant-coefficient or dy-
236 namic) equation for SGS turbulent kinetic energy [32] (FireFOAM, SIERRA/Fuego); or the
237 dynamic Smagorinsky model [33] (ISIS).
- 238 • For combustion: a global combustion equation combined with a closure expression for the
239 reaction rates based on either the Eddy Dissipation Model [34] (FDS, see also Ref. [31],
240 FireFOAM, ISIS) or a steady laminar flamelet model [35, 36] (SIERRA/Fuego).
- 241 • For radiation: a treatment based on either a solution of the radiative transfer equation (RTE)
242 combined with a closure expression for the emission term using a prescribed global radiative
243 loss fraction (FDS, see also Ref. [31], FireFOAM) or a solution of simplified equations based
244 on the $P1$ -approximation (ISIS) (note that results obtained with SIERRA/Fuego did not use
245 a radiation model).

246 These baseline models have a number of known limitations. First, the SGS turbulence models
247 are models that have been formulated for high-Reynolds-number, momentum-driven flow applica-
248 tions and that may not apply directly to fires which feature moderate-Reynolds-number, buoyancy-
249 driven flow and Rayleigh-Taylor instabilities. Second, the combustion models based on a global
250 combustion equation and the Eddy Dissipation Model are limited to configurations without ig-
251 nition/extinction phenomena and need to be modified to treat flame extinction in applications
252 to under-ventilated fires or suppressed fires. And third, the radiation models based on a pre-
253 scribed global radiative loss fraction are limited to configurations in which these fractions have
254 been measured and in which the fire regime does not deviate significantly from the conditions of
255 the measurements.

256 The identification of these limitations, the quantification of their relative importance, and ultimately their elimination through more advanced models are some of the objectives of the MaCFP
257 effort.
258

259 *2.2. Case 1: Turbulent Buoyant Plumes*

260 *2.2.1. Experiment*

261 The buoyant plume experiment selected for the first MaCFP workshop is a turbulent non-
262 reacting helium plume studied at a test facility called the Fire Laboratory for the Accreditation
263 of Models by Experimentation (FLAME) facility at Sandia National Laboratories (Sandia) [7, 37].
264 The original goal of the Sandia buoyant plume experiment was to provide comprehensive turbulent
265 flow velocity and species concentration statistics in a configuration that is representative of large-
266 scale pool fires without the complexities of chemical reactions and temperature variations [7]. The
267 1-m diameter source provides a plume in the fully-developed turbulent flow regime.

268 The 1-m diameter helium source was surrounded by a 0.51-m wide steel lip, representing the
269 injection plane and elevated 2.45 m above an annular ring which introduced a low-velocity co-flow
270 of ambient air [37]. The FLAME facility can be approximated as a 6.1-m cubic chamber covered
271 by a 2.4-m diameter extraction hood. Planar imaging measurements of velocity and species were
272 conducted using Particle Image Velocimetry and Planar Laser Induced Fluorescence, respectively.
273 Laser measurements were recorded at 200 Hz in a window approximately 0.86 m high and 1.2 m
274 wide, and providing an image of the near-field region (starting from the helium injection plane
275 and centered on the plume centerline). The measurement window includes near-field entrainment
276 zones on both sides of the plume; however, it does not include the lateral and vertical far-field. The
277 experimental uncertainty of the measured velocities and turbulent statistics are reported as 20%
278 and 30%, respectively. The uncertainty of the measured helium concentration is reported as 18%.
279 Inlet conditions are uniform to within 5% or less for the helium flow and within 10% for the air
280 coflow. The above uncertainties include run-to-run variability.

281 For the purpose of MaCFP, tests no. 25, 29, 32 and 36 were selected corresponding to repeat
282 runs with a helium inlet velocity of $0.339 \text{ m/s} \pm 1.3\%$, a flow Reynolds number $Re = 3194 \pm 0.6\%$,
283 a flow Richardson number $Ri = 69.53 \pm 6.5\%$ and a measured puffing frequency of 1.45 Hz [7].

284 *2.2.2. Simulations*

285 Three groups submitted computational results for Case 1: IRSN [38], NIST [39] and UGent [40].
286 Some of the information was also compared with past results published in the literature by Desjardin

287 *et al.* running an in-house solver [41]. IRSN used ISIS version 4.8.0 [21]; NIST used an official release
 288 of FDS (version 6.5.3) [18]; UGent used FireFOAM version 1.6 [20].

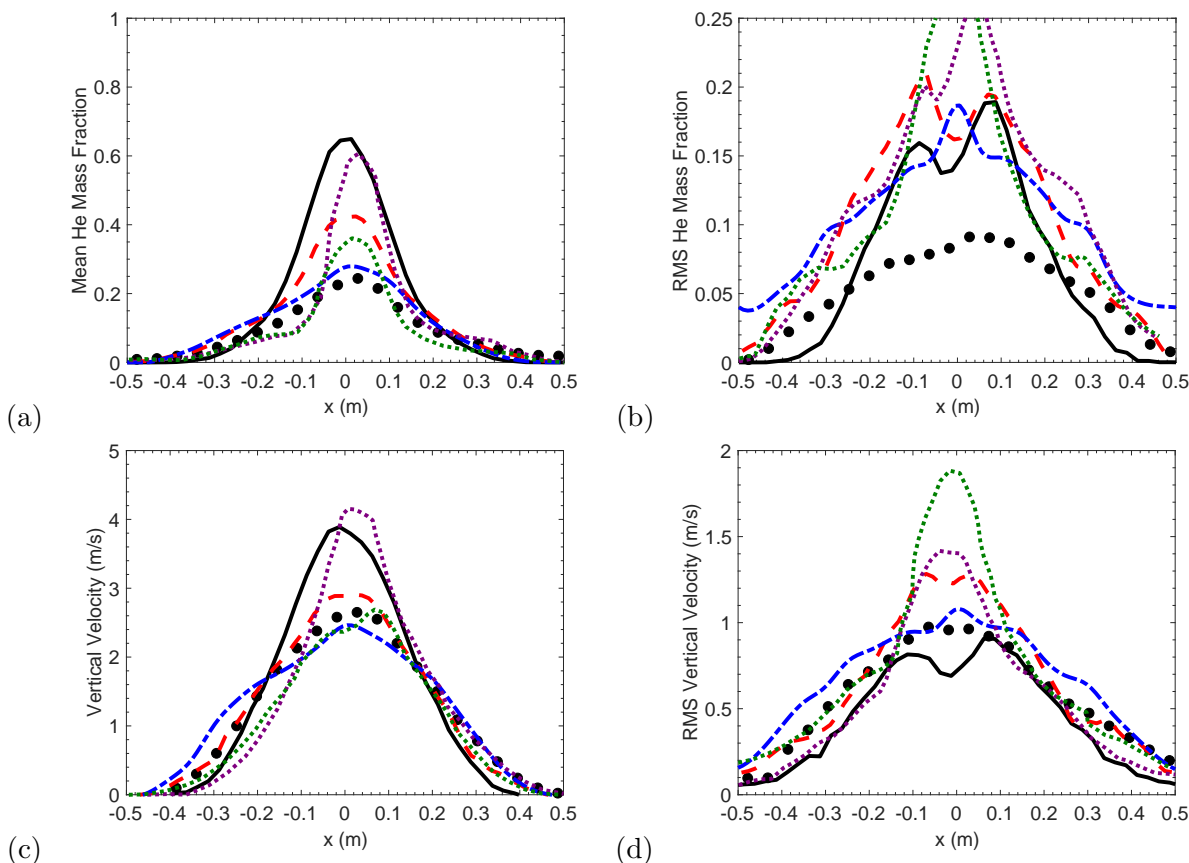


Figure 3: Case 1. Radial variations at $z = 0.4$ m of: (a) mean helium mass fraction; (b) *rms* helium mass fraction; (c) mean vertical velocity; (d) *rms* vertical velocity. Comparison between experimental data (black circles) and numerical results from IRSN (black solid line), NIST (red dashed line), UGent (blue dash-dotted line) and results from Ref. [41] (magenta and green dotted lines, corresponding to 5-cm and 3-cm resolution, respectively).

289 As discussed in section 2.1.2, the main question found in the design of a computational grid
 290 for LES simulations of the Sandia helium plume experiment is to decide whether to only require
 291 that the grid captures the large-scale dynamics occurring at length scale D or to also require
 292 that it captures the small-scale dynamics occurring in the intermittent boundary layer and the
 293 buoyancy-driven “thermals” at length scales δ_{BL} and $\delta_{thermals}$, respectively (see Figure 2). The
 294 former choice requires centimeter-scale resolution; the latter may require millimeter-scale resolution.
 295 The computational groups responded to this challenge in different ways: IRSN adopted a 2.5-
 296 cm resolution; NIST adopted a 1.5-cm resolution; UGent adopted a stretched grid with 1.23-cm
 297 resolution near the helium source. Note that previous work in Ref. [41] used both 3- and 5-cm

298 grid resolution. The computational domain in all simulations is much larger than the measurement
299 region (and is 3- or 4-m wide and 4-m high); however, it does not include all details of the full
300 facility.

301 Note that there were some variations among computational groups in the treatment of the
302 co-flow: the air co-flow velocity was introduced through an annular ring located below the helium
303 injection plane with a ring-level velocity approximately equal to 0.15-0.18 m/s [37]; IRSN did
304 attempt to model the annular ring but did not use the correct geometry and coflow velocity;
305 in contrast, NIST and UGent did not attempt to model the annular ring and assumed instead
306 simplified boundary conditions at the helium injection plane - NIST prescribed a small coflow
307 velocity equal to 0.01 m/s (at the injection plane) while UGent used free entrainment conditions.
308 Also, while NIST and UGent assumed an ambient pressure of 80.9 kPa (due to the high elevation
309 of the FLAME facility), IRSN incorrectly assumed an ambient pressure of 101.1 kPa.

310 Additional differences in the numerical treatment of the Sandia plume experiment include dif-
311 ferences in the choice of physical models (see section 2.1.3 for details on baseline choices). IRSN
312 used the baseline configuration of ISIS and NIST used the baseline configuration of FDS. UGent
313 deviated from baseline choices in FireFOAM and used the constant-coefficient Smagorinsky model
314 for subgrid-scale turbulence (additional information can be found in Ref. [42]).

315 The durations of the simulations and the durations over which numerical results were collected
316 and statistical moments were evaluated varied: IRSN, NIST and UGent chose to run their models
317 for 10 s, 20 s and 30 s, respectively (in Ref. [41], the model was run for 20 s); all groups except
318 IRSN chose to collect numerical results over the last 10 s, corresponding to approximately 14
319 puffing cycles; IRSN analyzed data over 3 s or approximately 4 puffing cycles (in this case, the
320 results should be analyzed with caution because the statistics may not be converged).

321 2.2.3. Summary

322 Figure 3 presents a representative sample of comparisons between measured and simulated
323 helium mass fractions and vertical flow velocity. The comparisons generally suggest that accuracy
324 increases with higher levels of grid resolution and that an accurate description of the flow statistics
325 in the near-field region requires a resolution of approximately 1 or 2 cm. Note that this corresponds
326 to 100 or 50 computational cells across the source diameter, a level of resolution that is much higher
327 than the usual requirement of providing a grid spacing that is 10-20 times smaller than the source
328 diameter D . In addition, even at this high level of resolution, the magnitude of the fluctuations

329 in helium mass fraction is not captured accurately (see Figure 3(b)). These inconsistent results
330 suggest that the dynamics associated with the presence of both thin boundary layers near the
331 edges of the helium source and small-scale “thermals” generated by secondary buoyant instabilities
332 play a significant role in determining near-field turbulent mixing properties. Note that the exact
333 impact of the inaccuracies in the near-field on the global flow features of the far-field have not been
334 characterized.

335 It is worth emphasizing that the experimental database describing the Sandia helium plume
336 experiment is of great value because it contains data on first and second-order statistical moments
337 of flow velocities and helium concentrations measured with high spatial resolution [7]. There are
338 also some limitations in the database that are worth pointing out for future studies: (1) the Sandia
339 database is limited to the plume near-field, *i.e.* to low elevations ($z \leq 1.2$ m, *i.e.* $z < (1.2 \times D)$),
340 and there is a need to provide similar data in the far-field; (2) the Sandia database does not
341 provide much information on the puffing cycle and there is a need to provide phase-averaged data
342 to characterize the coupling between large- and small-scale dynamics.

343 *2.3. Case 2: Turbulent Pool Fires with Gaseous Fuel*

344 *2.3.1. Experiments*

345 The gaseous pool fire experiments selected for the first MaCFP workshop correspond to a series
346 of natural gas flame experiments (Case 2a) studied at the National Institute of Standards and
347 Technology (NIST) [8] and a series of methane and hydrogen flame experiments (Case 2b) studied
348 at Sandia National Laboratories (Sandia) [9, 10]. The original goal of the NIST McCaffrey natural
349 gas flame experiment was to provide data to establish and/or validate engineering correlations for
350 mean temperature and mean vertical flow velocity along the center line of pool fires. The original
351 goal of the Sandia methane and hydrogen flame experiment was to provide comprehensive turbulent
352 flow velocity statistics in a configuration that is representative of large-scale pool fires.

353 The McCaffrey burner is a small-scale (0.3×0.3) m² square burner; in Ref. [8], the total heat
354 release rate was varied between 14.4 and 57.5 kW. **The NIST McCaffrey flames featured large-scale**
355 **unstable puffing motions at a frequency of 3 Hz.** Measurements of temperature and vertical flow
356 velocity were made using thermocouples and bi-directional probes. The Sandia burner is a 1-m-
357 diameter round burner located in a facility that can be approximated as a 6.1 m cube covered by
358 an extraction hood; in Refs. [9, 10], the total heat release rate was MW-scale (for the purpose of
359 MaCFP, tests no. 14, 24, 17 and 35 were selected corresponding to methane flames with a total

360 heat release rate equal to 1.59, 2.07 and 2.61 MW, and to a hydrogen flame with a total heat
361 release rate equal to 2.12 MW, respectively). The Sandia flames featured strong puffing motions **at**
362 **a frequency of 1.5 Hz** and the formation of thermals. Flow velocities were measured using Particle
363 Image Velocitometry; starting from the burner surface, measurements were made at high resolution
364 (with a spacing of 2 cm) and over a region approximately 0.9 m high and 1 m wide. Estimates of
365 errors in mean velocities for tests no. 14, 24, 17 and 35 range between 13% and 23%; estimates of
366 errors in *rms* velocities range between 13% and 28%.

367 2.3.2. Simulations

368 Four groups submitted computational results for Case 2a: FM Global [43], UGent [44], IRSN [45]
369 and NIST [46]. Four groups submitted computational results for Case 2b: UGent [47], NIST [48],
370 SNL [49] and UCantabria [50]. FM Global used a shared development version of FireFOAM
371 (FireFOAM-dev) [20]; UGent used FireFOAM version 2.4.x (Case 2a) and version 2.2.x (Case
372 2b) [20]; IRSN used ISIS version 4.8.0 [21]; NIST and UCantabria used an official release of FDS
373 (version 6.5.3) [18]; SNL used SIERRA/Fuego version 4.44 [22].

374 As discussed in section 2.1.2, the main **question** found in the design of a computational grid for
375 LES simulations of the NIST McCaffrey flame experiment **or the Sandia gas flame experiments is**
376 **to decide whether to only require that the grid captures the flow and flame features with length**
377 **scales comparable to the burner size or to also require that it captures the intermittent boundary**
378 **layer flame and the thermals that result from the puffing instability. For the NIST McCaffrey flame**
379 **experiment, all groups responded to this challenge in similar ways: FM Global and UGent adopted**
380 **a 1.25-cm resolution in the flame zone; IRSN adopted a 1-cm resolution; NIST adopted a 1.43-cm**
381 **resolution (the effective-pool-diameter-to-cell-size ratio varied between 24 and 34). For the Sandia**
382 **gas flame experiments, the computational groups responded to this challenge in different ways:**
383 **UGent and NIST adopted a 1.5-cm resolution in the flame zone; SNL presented results obtained**
384 **with a 2.5-cm and a 4-cm resolution; UCantabria adopted a 5-cm resolution (the pool-diameter-**
385 **to-cell-size ratio varied between 20 and 67). Note that while UGent and NIST decided to limit**
386 **the computational domain to a subset of the experimental facility, SNL and UCantabria chose to**
387 **include details of the full facility, including the co-flow arrangement and the extraction hood.**

388 Additional differences in the numerical treatment of the NIST McCaffrey and Sandia **flame**
389 experiments include differences in the choice of physical models (see section 2.1.3 for details on
390 baseline choices). For Case 2a, FM Global used the baseline configuration of FireFOAM except

391 for using a slightly simplified radiation treatment in which emission losses are correctly included
392 in the energy equation (using the global radiative loss fraction concept) but radiation transport
393 is ignored (*i.e.* the RTE equation is not solved) because comparisons to experimental data do not
394 require the evaluation of a heat flux at a remote surface; the values of the global radiative loss
395 fraction were prescribed using the measured values (varying between 17% and 27%). UGent also
396 used the baseline configuration of FireFOAM except for using the dynamic Smagorinsky model [33]
397 for subgrid-scale turbulence; the value of the global radiative loss fraction was prescribed as equal
398 to 20%; in the solution of the RTE, the discretization of angular space used 48 angles. IRSN used
399 the baseline configuration of ISIS. NIST used the baseline configuration of FDS: the values of the
400 global radiative loss fraction were prescribed using the measured values; in the solution of the RTE,
401 the discretization of angular space used 104 angles.

402 The durations of the simulations and the durations over which numerical results were collected
403 and statistical moments were evaluated varied: FM Global, IRSN, NIST and UGent chose to run
404 their models for 100 s, 11 s, 30 s and 50 s, and to collect numerical results over the last 80 s, 3 s,
405 20 s and 45 s (corresponding to approximately 240, 9, 60 and 135 puffing cycles), respectively. In
406 the simulation from IRSN, the results should be analyzed with caution because the statistics may
407 not be converged.

408 For Case 2b, UGent used the baseline configuration of FireFOAM except for using the constant-
409 coefficient Smagorinsky model for subgrid-scale turbulence and an emission/absorption treatment
410 of the RTE for radiation combined with a grey model (for methane flames, the global radiative loss
411 fraction was predicted to be equal to 24.8%); in the solution of the RTE, the discretization of angular
412 space used 48 angles (additional information can be found in Ref. [51]). NIST and UCantabria used
413 the baseline configuration of FDS (except that UCantabria used the Vreman model [31] for subgrid-
414 scale turbulence): the value of the global radiative loss fraction was prescribed as equal to 20%
415 (for methane flames) or 10% (for hydrogen flames); in the solution of the RTE, the discretization
416 of angular space used 104 angles. SNL used the baseline configuration of SIERRA/Fuego (but
417 without radiation).

418 The durations of the simulations and the durations over which numerical results were collected
419 and statistical moments were evaluated varied: NIST, SNL, UCantabria and UGent chose to run
420 their models for 20 s, 70 s, 50 s and 50 s, and to collect numerical results over the last 10 s, 30 s,
421 10 s and 45 s (corresponding to approximately 15, 45, 15 and 67 puffing cycles), respectively.

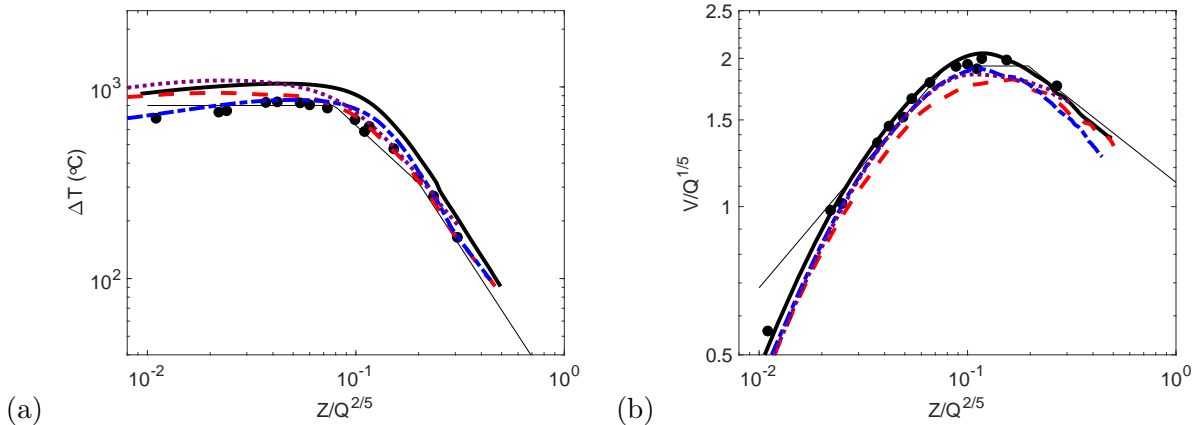


Figure 4: Case 2a. Vertical variations along the pool centerline (log-log plot): (a) mean excess temperature; (b) mean vertical velocity. Following standard scaling laws, vertical elevation is scaled by $Q^{2/5}$ while velocity is scaled by $Q^{1/5}$, with Q the total heat release rate. Comparison between experimental data (black circles), engineering correlations (thin black solid line, see Ref. [8]) and numerical results from FM Global (black solid line), IRSN (red dashed line), NIST (blue dash-dotted line), UGent (magenta dotted line). Case of a 33-kW flame.

422 2.3.3. Summary

423 For Case 2a, all simulations seem to correctly reproduce the gross features of the flame structure
 424 observed in the NIST McCaffrey flame experiment. Figure 4 presents a representative sample of
 425 comparisons between measured and simulated temperatures and vertical flow velocity. Note that
 426 in the NIST McCaffrey experiment, thermocouple measurements were not corrected for radiation
 427 losses and therefore should be interpreted with caution. NIST is the only computational group
 428 that used a thermocouple model to provide a sound basis for comparisons to the raw thermocouple
 429 measurements (the model is integrated inside the LES solver and uses the LES solution to simu-
 430 late deviations of thermocouple temperatures from gas temperatures [31]); other groups reported
 431 gas temperatures that require a correction before making a comparison to the raw thermocouple
 432 measurements; in Fig. 4, the temperatures reported by NIST are the simulated thermocouple tem-
 433 peratures, whereas the temperatures reported by other groups are the simulated gas temperatures.

434 It is worth emphasizing that while the experimental database describing the NIST McCaffrey
 435 natural gas flame experiment is a valuable starting point for model validation, there are, however,
 436 some obvious limitations in the database that are worth pointing out for future studies: (1) the
 437 database is limited to small-scale, weakly-to-moderately turbulent flames; and (2) the database is
 438 limited to temporal means and does not contain information on fluctuation magnitudes.

439 We now proceed to a discussion of Case 2b. All simulations seem to correctly reproduce the

440 gross features of the flame structure observed in the Sandia pool-like fire experiment. Figures 5-6
441 present a representative sample of comparisons between measured and simulated mean vertical and
442 radial velocities. Mean radial velocities are particularly important because they provide a measure
443 of the air entrainment process that determines the vertical mass flow rate in the flame and plume
444 regions (and thereby controls smoke production in fires): for instance, Figure 6 shows that at the
445 edge of the pool fire (*i.e.* at 0.5-m distance from the center of the burner), the radial velocity is
446 over-estimated by a factor close to two in the SNL and UCantabria simulations (at $z = 0.3$ m).
447 Additional comparisons can be found in [4]. Overall, the UGent and NIST simulations show good
448 agreement with experimental data and provide a satisfactory description of the flame structure.
449 The accuracy of the UCantabria simulation is limited by insufficient grid resolution.

450 It is worth emphasizing that the experimental database describing the Sandia methane and
451 hydrogen **gas flame** experiment is quite unique because it contains data on first and second-order
452 statistical moments of vertical/radial velocities measured with high spatial resolution [9, 10]. There
453 are also some limitations in the database that are worth pointing out for future studies: (1) the
454 Sandia database is limited to the flame near-field, *i.e.* to low elevations ($z \leq 0.9$ m, *i.e.* $z < D$),
455 and there is a need to provide data over the full flame region ($0 \leq z \leq L_f$); (2) the Sandia
456 database focuses on the flow structure in the flame region but does not contain information on the
457 temperature and radiation fields.

458 Finally, it is also worth noting that while research-level simulations may accept the compu-
459 tational cost associated with **high-resolution (*i.e.* with grids characterized by values of the pool-
460 diameter-to-cell-size ratio larger than 20)**, **engineering-level simulations will not accept that cost
461 and will use coarser grids (*i.e.* grids characterized by values of the pool-diameter-to-cell-size ratio
462 smaller than 10 or 20)**. These coarser-grid simulations require accurate subgrid-scale models: the
463 evaluation of current SGS models in simulations with representative engineering-level grids was not
464 part of the scope of the first MaCFP workshop and will be addressed in future editions.

465 2.4. Case 3: Turbulent Pool Fires with Liquid Fuel

466 2.4.1. Experiment

467 The liquid pool fire experiment selected for the first MaCFP workshop is a 0.305-m diameter
468 methanol pool fire previously studied at the University of Waterloo (UW) [11, 12]. The UW flame
469 was established over a modified liquid pan burner designed for free air entrainment. The burner
470 was operated at steady state with a gravity fuel feed of $1.35 \text{ cm}^3/\text{s}$ (1.07 g/s) for a total heat release

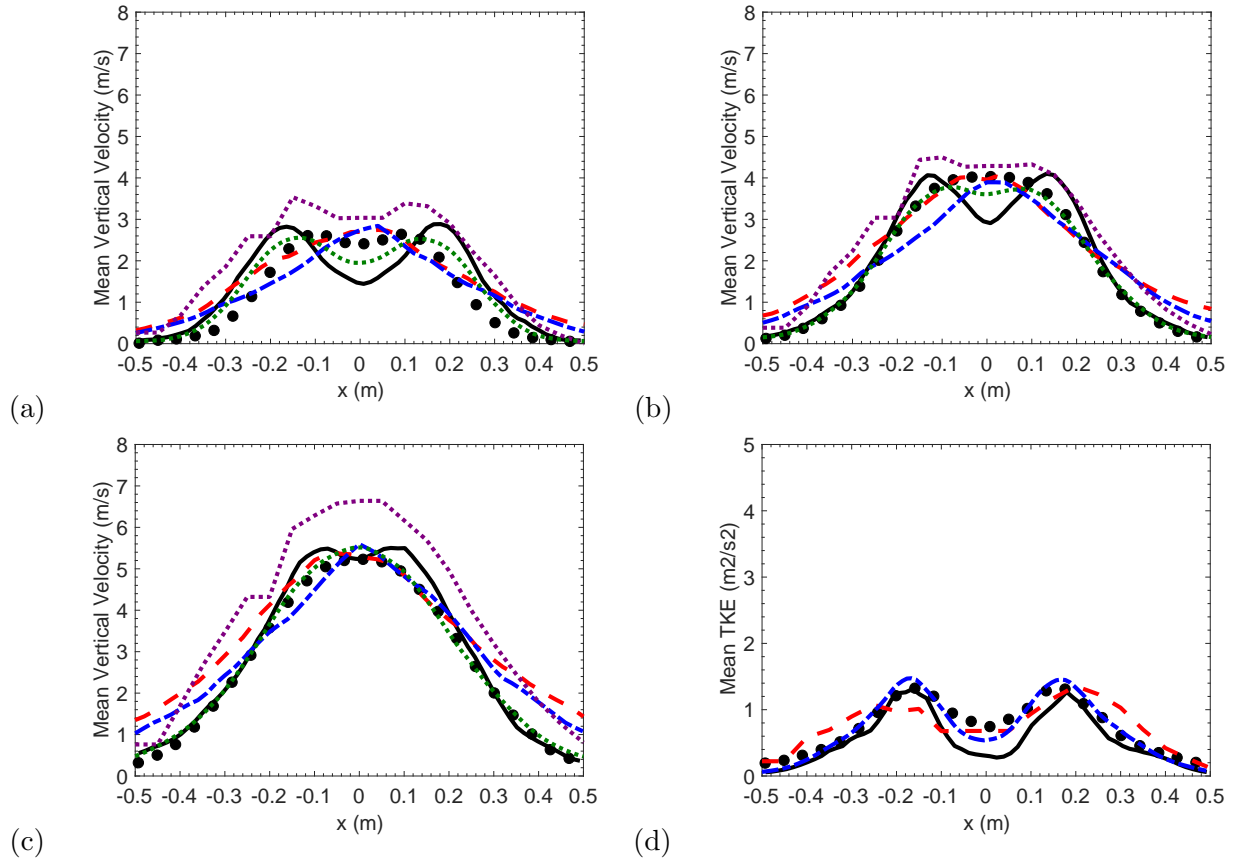


Figure 5: Case 2b. Radial variations of mean vertical velocity at elevation: (a) $z = 0.3$ m; (b) $z = 0.5$ m; (c) $z = 0.9$ m; and (d) radial variations of mean (resolved) turbulent kinetic energy at $z = 0.5$ m. Comparison between experimental data (black circles) and numerical results from NIST (black solid line), SNL (red dashed and blue dash-dotted lines); UCantabria (magenta dotted line); UGent (green dotted line). Case of a 2.07-MW methane flame (test 24).

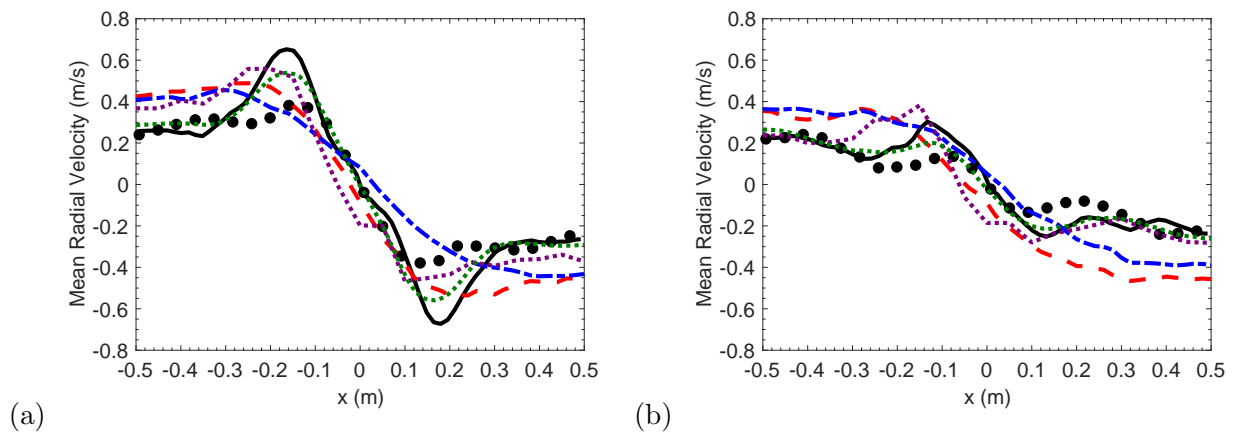


Figure 6: Case 2b. Radial variations of mean radial velocity at elevation: (a) $z = 0.3$ m; (b) $z = 0.5$ m. See caption of Fig. 5.

471 rate of 22.6 kW. The height of the burner lip above the liquid fuel surface was 1 cm. The flame
472 was approximately 0.5-m-high and featured a strong puffing instability; the frequency of oscillation
473 was 2.8 Hz.

474 Time-resolved velocity (using two component, forward-scatter Laser Doppler Anemometry) and
475 temperature (using 50 micron diameter, bare-wire Pt-Pt-10%Rh thermocouples with 75-100 micron
476 beads) were measured in the highly-fluctuating region of the flame, *i.e.* up to radial positions located
477 16 cm from the pool fire centerline and up to 30 cm vertical elevation. Direct and Schlieren pho-
478 tography of the luminous flame were used to characterize the macroscopic and oscillatory behavior
479 of the flame.

480 Time series of data were averaged to provide mean and *rms* values as well as correlation coef-
481 ficients [12]. Errors in mean and *rms* velocities and mean temperatures were estimated as $\pm 5\%$ at
482 95% confidence; errors in Reynolds stresses were estimated as $\pm 15\%$ at 95% confidence [11]. Errors
483 in *rms* temperatures and velocity-temperature correlations were difficult to estimate and were not
484 quantified. No correction was made to temperature measurements for radiation or catalytic effects
485 (estimated to be less than 5%).

486 2.4.2. Simulations

487 Three groups submitted computational results for Case 3: UGent [52], UMD [53] and VTT [54].
488 UGent used FireFOAM version 2.2.x [20]; UMD used a shared development version of FireFOAM
489 (FireFOAM-dev) [20]; VTT used an official release of FDS (version 6.5.3) [18].

490 Two specific **questions** were found in the design of a computational grid for LES simulations
491 of the UW pool fire experiment. The first **question** is to **decide whether the grid should capture**
492 the intermittent boundary layer flame and thermals that result from the puffing instability. As
493 discussed in section 2.1.2, this **may require** millimeter-scale resolution. The second challenge is the
494 presence of a 1-cm-high pool lip in the UW experiment. The presence of the lip leads to complex
495 flow patterns close to the edges of the methanol pool surface that require millimeter-scale resolution
496 to be correctly captured by the computational grid.

497 The computational groups responded to these two challenges in different ways: UGent adopted
498 a 5-mm resolution in the flame zone (both in the horizontal and vertical directions) and included
499 the burner lip in the numerical configuration; UMD adopted a 2.5-mm resolution in the vertical
500 direction but also used a coarser 10-mm resolution in the horizontal directions and did not account
501 for the presence of the lip; VTT adopted a 2.5-mm resolution (both in the horizontal and vertical

502 directions) and did not account for the presence of the lip. The pool-diameter-to-cell-size ratio
 503 varied between 30 and 120.

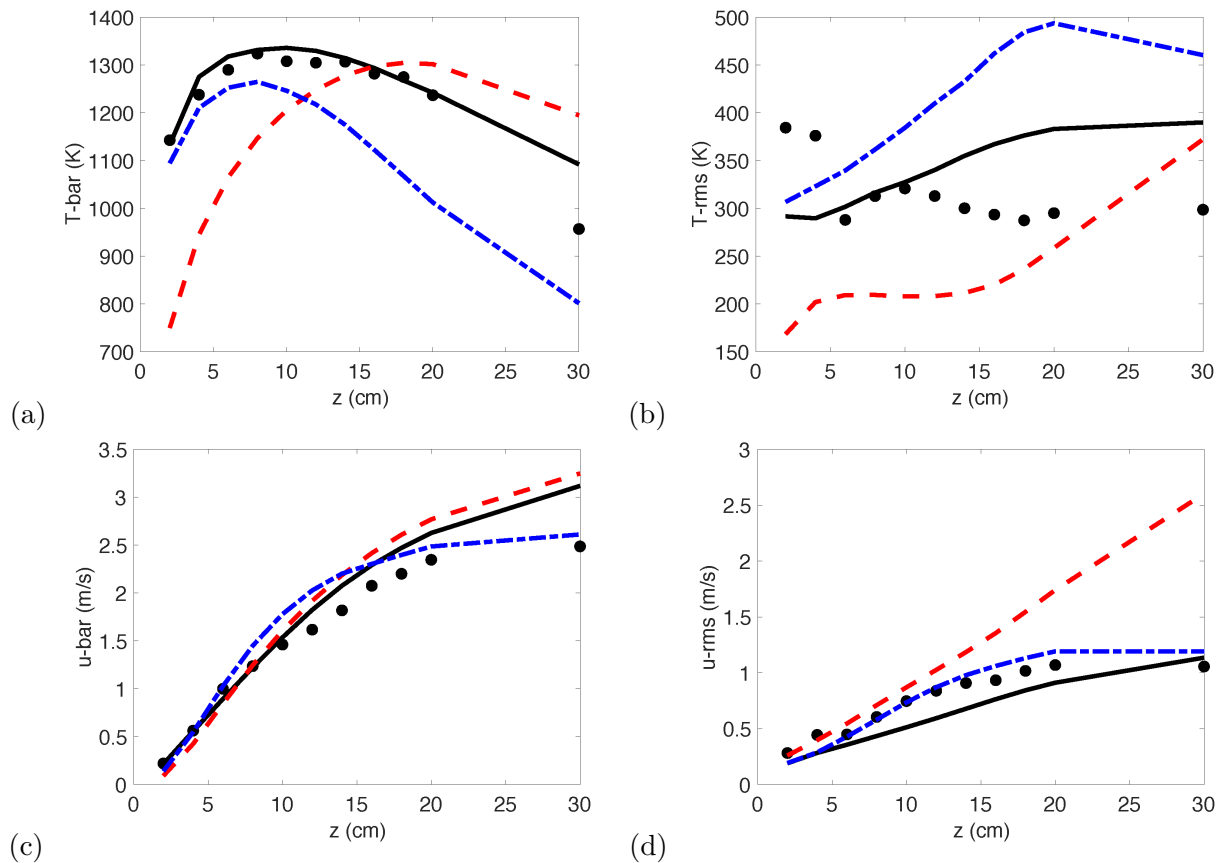


Figure 7: Case 3. Vertical variations along the pool centerline: (a) mean temperature; (b) rms temperature; (c) mean vertical velocity; (d) rms vertical velocity. Comparison between experimental data (black circles) and numerical results from UGent (black solid line), UMD (red dashed line), VTT (blue dash-dotted line).

504 An important difference in the numerical treatment of the UW experiment is that while UGent
 505 and UMD prescribed the fuel evaporation rate using the measured mean experimental value (1.07
 506 g/s), VTT adopted a more ambitious treatment in which the fuel evaporation rate is calculated
 507 as a function of the gas-to-liquid thermal feedback. In the simulation performed by VTT, the
 508 fuel evaporation rate is under-predicted by a factor close to 1.7 leading to a flame size of 13 kW
 509 (compared to 22 kW in simulations by UGent and UMD).

510 Additional differences in the numerical treatment of the UW experiment include differences in
 511 the choice of physical models (see section 2.1.3 for details on baseline choices). UGent deviated
 512 from baseline choices in FireFOAM and used the dynamic Smagorinsky model [33] for subgrid-
 513 scale turbulence, the Eddy Dissipation Concept (EDC) model [55] for combustion, and an emis-

514 sion/absorption treatment of the RTE for radiation combined with a Weighted-Sum-of-Gray-Gases
 515 model [56] for gas radiation (the global radiative loss fraction was predicted to be equal to 16.4%, a
 516 value that is close to the empirically-determined value of 17-18%); in the solution of the RTE, the
 517 discretization of angular space used 72 angles. UMD used the baseline configuration of FireFOAM;
 518 the value of the global radiative loss fraction was prescribed as equal to 18%; in the solution of the
 519 RTE, the discretization of angular space used 16 angles. VTT used the baseline configuration of
 520 FDS: the values of the global radiative loss fraction was prescribed as equal to 17%; in the solution
 521 of the RTE, the discretization of angular space used 104 angles.

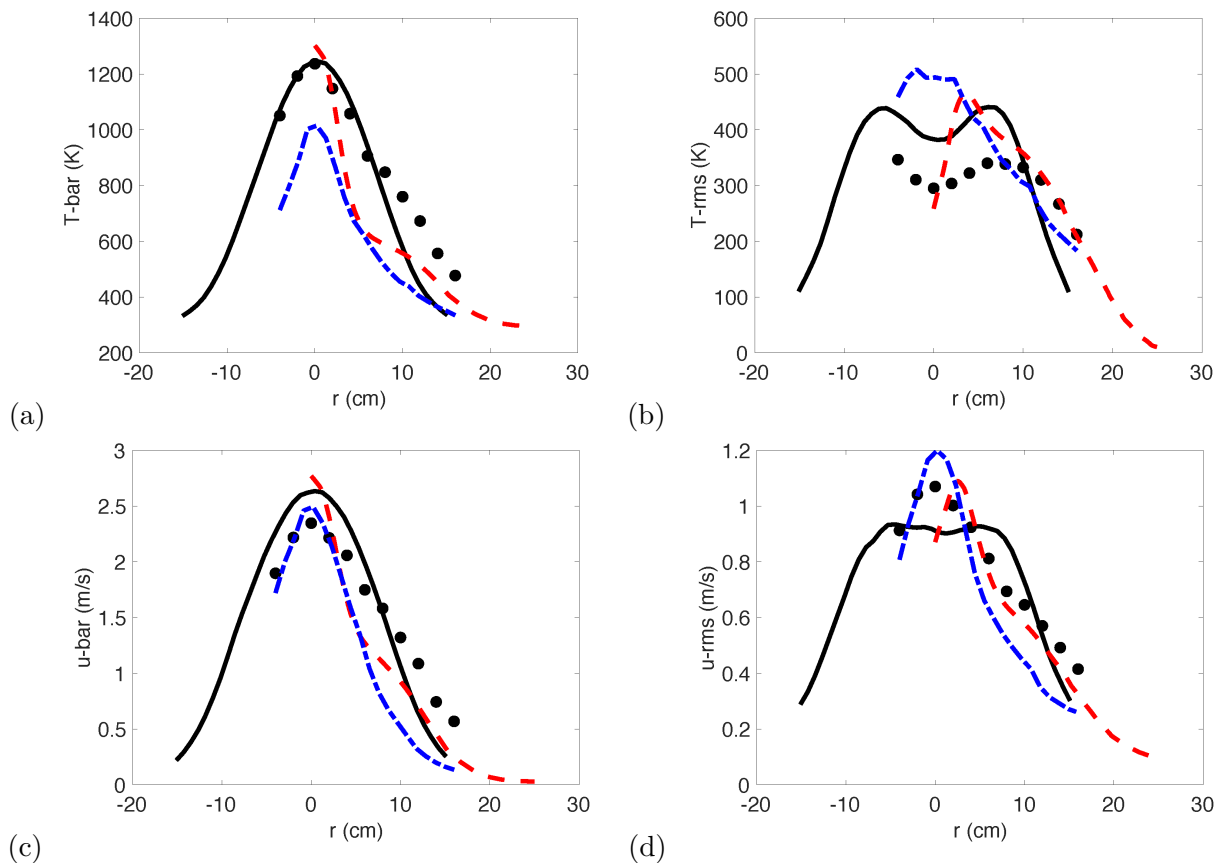


Figure 8: Case 3. Radial variations at elevation $z = 20$ cm: (a) mean temperature; (b) *rms* temperature; (c) mean vertical velocity; (d) *rms* vertical velocity. See caption of Fig. 7.

522 The durations of the simulations and the durations over which numerical results were collected
 523 and statistical moments were evaluated varied: UGent, UMD and VTT chose to run their models
 524 for 65 s, 60 s and 15 s, and to collect numerical results over the last 60 s, 50 s and 10 s (corresponding
 525 to approximately 168, 140 and 28 puffing cycles), respectively.

526 2.4.3. Summary

527 All simulations correctly reproduce a pulsating flame with a frequency of oscillation close to
528 the measured value (2.8 Hz): 2.8 Hz (UGent), 2.2 Hz (UMD), 3 Hz (VTT). Figures 7 and 8
529 present a small representative sample of comparisons between experimental data and numerical
530 simulations. Additional comparisons can be found in [4]. Overall, the UGent simulation shows good
531 agreement with experimental data and provides a satisfactory description of the flame structure.
532 The accuracy of the UMD simulation is limited by insufficient grid resolution. The accuracy of the
533 VTT simulation is limited by an inaccurate prediction of the total heat release rate.

534 It is worth emphasizing that the experimental database describing the UW methanol pool fire
535 experiment is quite unique because it not only contains data on first and second-order statistical
536 moments of temperature and vertical/radial velocities, but also contains data on Reynolds shear
537 stresses and turbulent heat fluxes [11, 12]. There are also some limitations in the database that
538 are worth pointing out for future studies: (1) the UW database is limited to the flame near-field,
539 *i.e.* to low elevations ($z \leq 30$ cm), and there is a need to provide data over the full flame region
540 ($0 \leq z \leq L_f$, where L_f is the flame height; $L_f \approx 0.5$ m in the UW experiment); (2) the flame is
541 only weakly turbulent and there is a need to provide data for larger flame sizes, *i.e.* for larger pool
542 diameters ($D \geq 1$ m); (3) the thermal feedback is not characterized and there is a need to provide
543 data on the convective/radiative heat flux at the liquid fuel surface.

544 2.5. Case 4: Turbulent Wall Fires

545 2.5.1. Experiment

546 Turbulent fire on a vertical surface is a canonical configuration representing the upward flame
547 spread problem, typical in many practical fire scenarios. The FM Global vertical wall flame ex-
548 periment selected for the first MaCFP workshop [13, 14] is a series of meter-scale wall fires (the
549 total heat release rate is several 100s of kW) realized by an array of vertically-stacked water-cooled
550 porous gas burners with prescribed fuel supply rates. The setup conveniently decouples the gas-
551 phase fire dynamics and corresponding heat transfer from the solid-phase pyrolysis, and thereby
552 achieves a statistically steady-state condition ideal for experimental measurements and CFD model
553 validation. The original goal of the experiment was to provide data to establish theoretical models
554 and correlations for radiative and convective gas-to-solid heat transfer in wall fires. The fuel type
555 and the fuel injection rate were varied in the tests. The total flame-to-wall heat flux as well as
556 inward and outward flame radiation intensities were measured at different elevations. Other mea-

557 surements included wall-normal profiles of gas temperature and flow velocity and vertical profiles of
558 the soot depth. This experimental configuration provides MaCFP with the following: a canonical
559 configuration that brings data on flame-wall interactions with realistic scales and buoyant turbu-
560 lent flow conditions; a simplified statistically-stationary configuration with different forms of heat
561 transfer; and decoupled solid- and gas-phase processes through controlled fuel injection.

562 2.5.2. Simulations

563 Two groups submitted computational results for Case 4: FM Global [57] and NIST [58]. FM
564 Global used FireFOAM version 2.2.x [20]; NIST used an official release of FDS (version 6.5.3) [18].

565 As discussed in section 2.1.2, the main challenge found in the design of a computational grid for
566 LES simulations of the FM Global vertical wall flame experiment is to provide suitable grid resolu-
567 tion to capture the thin turbulent boundary layer flame. This requires millimeter-scale resolution.
568 FM Global and NIST responded to this challenge in a similar way and adopted a 3-mm resolution
569 in the near-wall flame region. In a previous study of the same wall fire configuration [59], this level
570 of spatial resolution was found to be adequate for grid-resolved LES simulations (*i.e.* for simula-
571 tions that capture the wall gradients and are performed without using wall models). NIST chose
572 to apply the 3-mm resolution in all directions and across the entire computational domain. FM
573 Global chose a lower-cost computational grid and applied the 3-mm resolution in the wall-normal
574 direction while using a 7.5-mm resolution in the spanwise and vertical directions (parallel to the
575 wall) and also used a coarser mesh in the far field.

576 Additional differences in the numerical treatment of the wall flame experiment include differ-
577 ences in the choice of physical models (see section 2.1.3 for details on baseline choices). FM Global
578 used the baseline configuration of FireFOAM except for using the WALE model [60] for subgrid-
579 scale turbulence and for correct behavior in the near-wall region; the values of the global radiative
580 loss fraction were prescribed using the measured values (to account for the radiation absorption
581 from the fuel and cold soot in the near wall region, the prescribed values were chosen as 75% of
582 the values of the radiative loss fraction measured in corresponding wall-free configurations); in the
583 solution of the RTE, the discretization of angular space used 16 angles. NIST used the baseline
584 configuration of FDS except for using an emission/absorption treatment of the RTE based on a
585 simplified soot formation model (using a prescribed soot yield) and a grey model for soot radiation
586 as well as a wide-band (6 bands) model for gas radiation (with coefficients calibrated by the narrow-
587 band model called RadCal [61]); in the solution of the RTE, the discretization of angular space

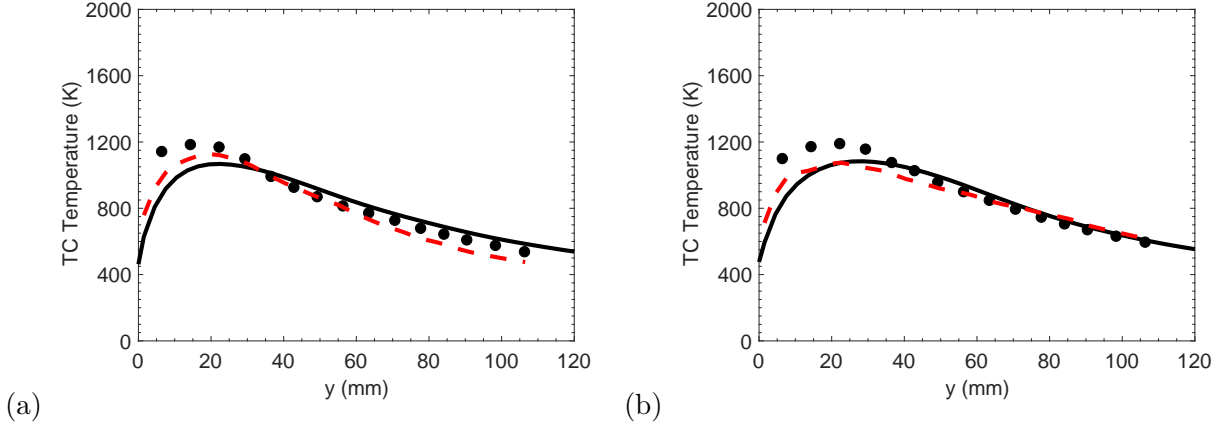


Figure 9: Case 4. Wall-normal variations of mean thermocouple temperature at $z = 0.77$ cm and for a fuel supply rate equal to: (a) $12.68 \text{ g/m}^2/\text{s}$; (b) $17.05 \text{ g/m}^2/\text{s}$. Comparison between experimental data (black circles) and numerical results from FM Global (black solid line) and NIST (red dashed line). Case of a propylene flame.

588 used 104 angles. Note that in its baseline configuration, FDS uses a Smagorinsky model with Van
 589 Driest wall functions [31] to estimate a turbulent viscosity at the wall and a Nusselt-number-based
 590 convective heat transfer model to estimate the convective heat flux at the wall [31].

591 Interestingly, FM Global and NIST differ significantly in their near-wall treatment: while FM
 592 Global follows the modeling choices of Ref. [59] and adopts a wall-resolved approach in which the
 593 wall convective heat flux is calculated by direct differentiation of the resolved temperature field
 594 (no wall model is used), NIST adopts a wall-modeled approach in which the wall convective heat
 595 flux is reconstructed by wall functions and Nusselt-number correlations. It is not clear whether
 596 the wall-modeled approach adopted by NIST converges towards a wall-resolved approach in case
 597 of sufficient grid resolution.

598 2.5.3. Summary

599 We limit our discussion to the case of propylene fuel (additional results can be found in [4]).
 600 Both FM Global and NIST simulations qualitatively reproduce the variations of the wall heat flux
 601 with vertical elevation as well as the variations of the wall heat flux in response to changes in
 602 the fuel supply rate. Figure 9 presents a representative sample of comparisons between measured
 603 and simulated thermocouple temperatures (note that in these comparisons, both FM Global and
 604 NIST simulations use a thermocouple model that is integrated inside the solvers and that uses
 605 the LES solution to simulate the deviations of thermocouple temperatures from gas temperatures).
 606 Figure 10 presents a sample of comparisons between measured and simulated wall heat fluxes.
 607 The NIST simulations overpredict the total heat flux by approximately 50% in most scenarios; in

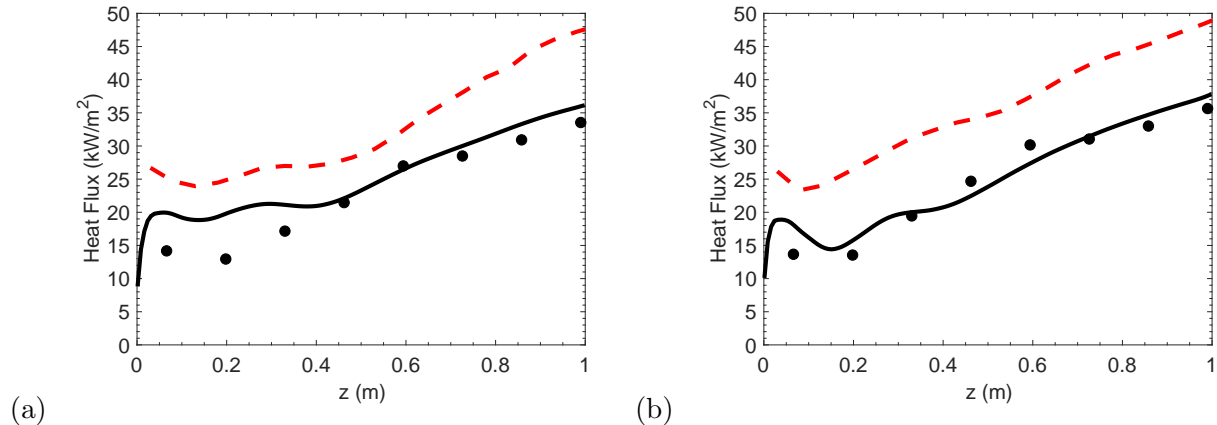


Figure 10: Case 4. Vertical variations of the mean total wall heat flux for a fuel supply rate equal to: (a) 12.68 $\text{g}/\text{m}^2/\text{s}$; (b) 17.05 $\text{g}/\text{m}^2/\text{s}$. See caption of Fig. 9.

608 contrast, the FM Global simulations show good agreement with experimental data. These results
 609 may be explained by the choice made in the FM Global simulations to use a semi-empirical radiation
 610 model with an experimentally-determined global radiative loss fraction compared to the choice made
 611 in the NIST simulations to use a more fundamental description through an emission/absorption
 612 treatment of the RTE. Note that because the experimental database does not include information on
 613 the convective and radiative components of the wall heat flux, this information was not extracted
 614 from the simulations. Future simulations of this case should analyze these components (see for
 615 instance Ref. [59]) and also bring information on the relative weight of soot radiation and gas
 616 radiation.

617 In closing, the experimental database describing the FM Global vertical wall flame experiment is
 618 quite unique because it brings fundamental information on gas-to-solid heat transfer processes that
 619 are a controlling factor in flame spread problems. There are also some limitations in the database
 620 that are worth pointing out for future studies: (1) the database is limited to temporal means and
 621 does not contain information on fluctuation magnitudes; (2) the thermal feedback is characterized
 622 in terms of a total wall heat flux but does not contain information on the convective and radiative
 623 components of the wall heat flux, nor on the soot and gas contributions to the radiative component
 624 of the wall heat flux.

625 Furthermore, as already pointed out in section 2.3, it is worth noting that while research-
 626 level simulations may accept the computational cost associated with millimeter-scale resolution,
 627 engineering-level simulations will not accept that cost and will use coarser grids that require wall
 628 models. The development and evaluation of these models is part of the future objectives of MaCFP.

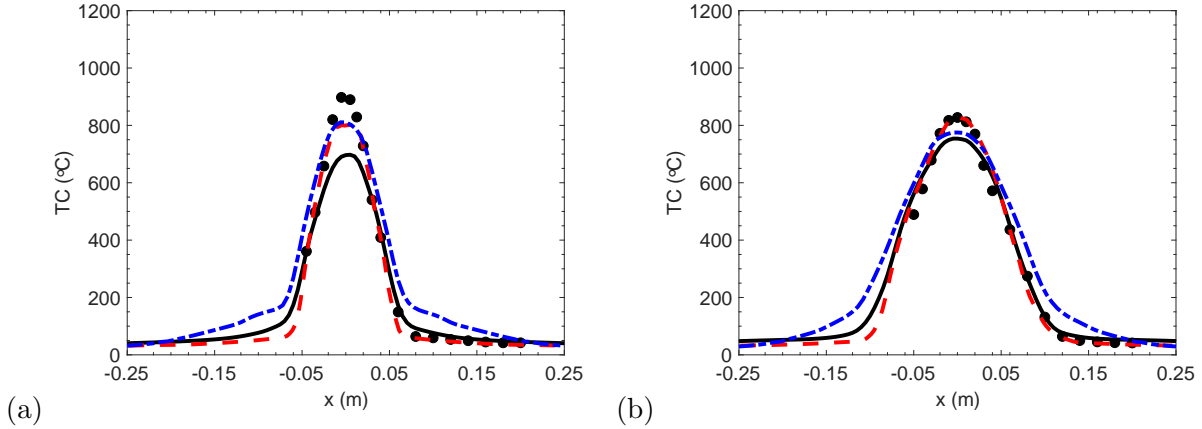


Figure 11: Case 5. Cross-flame variations of mean thermocouple temperature at: (a) $z = 0.125$ m; (b) $z = 0.25$ m. Comparison between experimental data (black circles) and numerical results from FM Global (black solid line), NIST (red dashed line), UMD (blue dash-dotted line). Case of a methane flame with $X_{O_2} = 18\%$.

629 2.6. Case 5: Flame Extinction

630 2.6.1. Experiment

631 The flame extinction experiment selected for the first MaCFP workshop is a canonical line-fire
 632 configuration with controlled co-flow studied at the University of Maryland (UMD) [15–17]. The
 633 UMD turbulent line burner facility allows the study of a buoyancy-driven, turbulent diffusion flame
 634 exposed to environments of decreasing oxygen strength, down to the oxygen extinction limit, and
 635 thereby provides fundamental information relevant to fire suppression due to under-ventilation or
 636 due to the activation of an inert gas system. The facility comprises a sand-filled, stainless-steel fuel
 637 port, slot burner, 5-cm-wide and 50-cm-long. Controlled suppression of the flame is achieved via
 638 the introduction of nitrogen gas into the co-flowing oxidizer stream.

639 Both methane and propane fuels were utilized. Assuming complete combustion, the total heat-
 640 release rate was 50 kW for both fuels (the flame was approximately 0.5 m-high). The quantitative
 641 metric of suppression is the global combustion efficiency, η , reported as a function of the coflow
 642 oxygen strength and measured using oxygen consumption and carbon dioxide generation calorime-
 643 try [17]. The mole fraction of oxygen, X_{O_2} , was measured using a paramagnetic oxygen analyzer
 644 via a probe located inside the oxidizer port. Infrared radiative emissions were measured using a
 645 water-cooled Schmidt-Boelter heat-flux transducer; heat flux data were then converted to global
 646 radiative loss fractions using a weighted multipoint radiation source model [16]. Visible flame height
 647 was measured using a video camera. A limited set of temperature measurements is also available
 648 for methane fuel and $X_{O_2} = 18\%$.

649 *2.6.2. Simulations*

650 Three groups submitted computational results for Case 5: FM Global [62], NIST [63] and
651 UMD [64]. FM Global and UMD used a shared development version of FireFOAM (FireFOAM-
652 dev) [20]; NIST used an official release of FDS (version 6.5.3) [18].

653 As discussed in section 2.1.2, the main challenge found in the design of a computational grid for
654 LES simulations of the UMD turbulent line flame experiment is to provide suitable grid resolution
655 to capture the controlling length scale of the burner, *i.e.* the burner width (5 cm). This requires
656 millimeter-scale resolution. The computational groups responded to this challenge in a similar way
657 and adopted a resolution of 5-mm (FM Global), 3.125-mm (NIST) and 4.2 mm (UMD) in the flame
658 region.

659 Additional differences in the numerical treatment of the line flame experiment include differences
660 in the choice of physical models (see section 2.1.3 for details on baseline choices). FM Global and
661 UMD used the baseline configuration of FireFOAM except for the addition of a flame extinction
662 model based on the concept of a critical Damköhler number for premixed eddies [65] (FM Global)
663 or the concept of a critical Damköhler number for diffusion flames [66] (UMD); the values of the
664 global radiative loss fraction were prescribed using the measured values [16]; in the solution of the
665 RTE, the discretization of angular space used 16 angles. NIST used the baseline configuration of
666 FDS except for the addition of a flame extinction model based on the concept of a critical flame
667 temperature [67, 68]; in the solution of the RTE, the discretization of angular space used 700 angles
668 (the large number of angles is due to the fact that the NIST simulation included the heat flux gauge
669 located at 1-m distance from the flame and was motivated by the desire to avoid any potential ray
670 effect).

671 *2.6.3. Summary*

672 All simulations seem to reproduce the overall structure of the turbulent flame (see Fig. ??).
673 Figure 11 presents comparisons between measured and simulated thermocouple temperatures per-
674 formed at quarter-flame height and at mid-flame height (note that in these comparisons, all simu-
675 lations use a thermocouple model). Figure 11 suggests that the level of agreement between exper-
676 imental data and numerical results is encouraging; discrepancies are however observed, especially
677 at quarter-flame height, and those may be attributed to inaccuracies in the combustion model, in
678 the thermal radiation model, or in the coupling of these models.

679 Figure 12 presents comparisons between measured and simulated combustion efficiencies as a

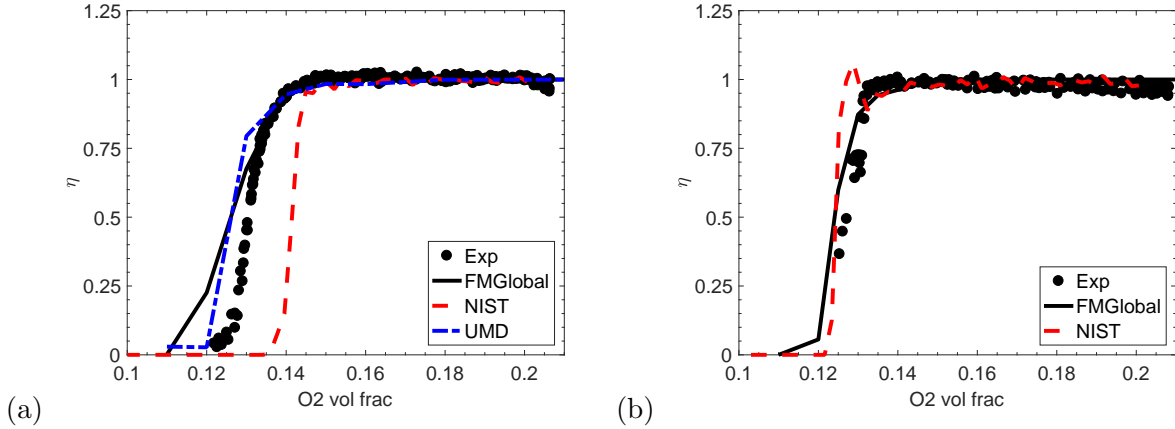


Figure 12: Case 5. Variations of the global combustion efficiency with the coflow oxygen mole fraction. (a) methane flame; (b) propane flame. See caption of Fig. 11.

680 function of the coflow oxygen strength, for both methane and propane flames. All simulations
 681 correctly reproduce the binary nature of the flame response: the combustion efficiency remains
 682 close to 1 for X_{O_2} above the extinction limit and abruptly decreases to 0 at this limit (*i.e.* for
 683 $X_{O_2} \approx 12\text{-}14\%$). The exact value of the oxygen extinction limit is predicted within $\pm 10\text{-}20\%$.
 684 While these results are encouraging, it is worth emphasizing that the flame extinction models are
 685 complex (they in fact rely on a description of both extinction and re-ignition phenomena [65, 66, 68])
 686 and that the models used in the FM Global, NIST and UMD simulations are based on different
 687 representations of the physics. Thus, the UMD turbulent line flame database is not capable of
 688 differentiating between the three flame extinction models and therefore does not provide sufficient
 689 insight into the underlying physics of flame suppression. Also, as mentioned in section 2.1.2, it
 690 is important to recognize that the FM Global and UMD flame extinction models (and to a lesser
 691 extent the NIST flame extinction model) were originally tuned against data obtained from the same
 692 UMD experiment. Therefore, the simulations should be interpreted as *calibration* tests rather than
 693 validation tests.

694 Note that the UMD turbulent line flame database has been recently enhanced with new micro-
 695 thermocouple measurements and has also been extended to the case of flame suppression by a
 696 water mist [15]. These new developments should be incorporated into MaCFP. The addition of
 697 micro-thermocouple measurements will provide much needed data to characterize the details of the
 698 flame structure (and will provide both first- and second-order statistics). More information on flow
 699 velocity as well as on gas and soot radiation will also be needed in order to unravel the respective
 700 effects of combustion and thermal radiation.

701 *2.7. Current and Future Plans*

702 The gas phase session of the June-2017 MaCFP workshop provided a first opportunity to demon-
703 strate the benefits and potential impact of activities organized by the IAFSS MaCFP Working
704 Group. The session provided a community-wide forum for in-depth technical discussions of a first
705 suite of experimental-computational comparisons corresponding to an initial list of target experi-
706 ments. The session was well attended (with 120 registered participants) and the first general lesson
707 from the workshop is that MaCFP successfully responds to a need for greater levels of integration
708 and coordination in fire research. The fire science community is small, fragmented and geograph-
709 ically dispersed: MaCFP is an effort to meet the resulting organizational challenge, to build an
710 international collaborative framework, and to provide a critical mass of researchers for topics cen-
711 tral to the development of a fundamental understanding of fire phenomena. While MaCFP is
712 currently focused on building a collaborative framework between computational and experimental
713 researchers around the topic of the validation of CFD-based fire models, we envision that MaCFP
714 can be extended to incorporate efforts focused on other topics, or can be emulated and inspire other
715 efforts.

716 The gas phase session of the first MaCFP workshop also led to a number of technical lessons and
717 outcomes that will help shape the future activities of MaCFP. First, as discussed in section 2.1, the
718 performance of CFD-based fire models depends on both the quality of the computational grid (and
719 in particular its ability to resolve the dynamically-controlling length scales of the simulated problem)
720 and the accuracy of the physical models (used to describe subgrid-scale turbulence, combustion,
721 radiation, *etc*). In this context, it is important to emphasize that the submissions made by the
722 seven different computational groups represented at the MaCFP workshop correspond to fine grid
723 resolution at the millimeter- or centimeter-scale. Under high-resolution simulation conditions, the
724 impact of numerical errors is reduced and many of the discrepancies between experimental data
725 and computational results may be attributed to modeling errors. While fine-grained simulations
726 are considered as a necessary step, and provide valuable insights into the accuracy of physical
727 models, they are not representative of engineering-level simulations that typically use coarser grids:
728 there is therefore an unmet need for MaCFP to also evaluate physical models in coarse-grained
729 simulations that are more representative of the CFD practice. This will be addressed in future
730 editions of the MaCFP workshop series. Note that one objective of the MaCFP Working Group
731 is to develop guidelines for CFD practitioners for the design of the computational grid as well as
732 reference material on the domain of validity of the different physical models available in current

733 CFD-based fire models.

734 Discussions of the different submissions by computational groups revealed a number of limita-
735 tions in the submitted results that are worth pointing out for the planning of future workshops: (1)
736 the comparisons between different computational results were of diminished value because there
737 was no required specific levels of spatial resolution (for the flow solver) or required grid convergence
738 studies; (2) the same comparisons were also of diminished value because there was no required spe-
739 cific levels of angular resolution (for the radiation solver) or required angular convergence studies;
740 (3) the presentations of computational results for different cases obtained by the same modeling
741 group were of diminished value because there was no requirement to define a baseline model con-
742 figuration that would be applied to all simulated cases considered by a particular group and no
743 requirement to provide a justification for possible variations in modeling choices; (4) the compar-
744 ison between experimental data and simulation results could be improved by specifying a scheme
745 or a metric to quantify discrepancies in comparisons of experimental data and simulation results;
746 and (5) the comparison between experimental data and simulation results could be improved by
747 including experimental uncertainties in the comparative plots.

748 Furthermore, discussions of the different target experiments revealed a number of limitations in
749 available experimental databases that are worth pointing out for future studies: (1) the databases
750 are often limited to small-scale, weakly-to-moderately turbulent flames and there is a need to
751 provide more data for large-scale fully-developed turbulent flames; (2) the databases are often
752 limited to measuring temporal means and there is a need to provide data on fluctuation magnitudes;
753 (3) the databases are often limited to the flame near-field and there is a need to provide data over
754 the full flame region; and (4) the databases are often focused on characterizing the flow field or
755 the temperature field, but not both, and there is a need to provide more comprehensive data sets
756 including flow velocities, temperatures, and also soot volume fractions, radiation intensities and
757 heat fluxes to surfaces. Note that the availability of quality data on radiation intensities and heat
758 fluxes to surfaces is a requirement for future progress on simulations of flame spread phenomena.
759 We hope that the wish list above will inspire a new generation of experimentalists and motivate
760 new experimental studies.

761 Finally, as discussed in section 1, the initial list of target experiments selected for the first
762 MaCFP workshop had a limited scope corresponding to (mostly) non-sooting or only weakly-sooting
763 flames, supplied with gaseous or liquid fuel, and without compartment effects. There is now a need
764 to extend the scope of MaCFP to include target experiments that bring detailed information on a

765 number of key fire processes, for instance: thermal radiation, soot formation and oxidation, flame
766 spread along solid flammable materials (see the next section), and also ignition phenomena and
767 compartment effects. In addition, the application of current fire models to the simulation of water-
768 based fire suppression systems (*i.e.* sprinkler or mist systems) requires additional experimental data
769 and validation tests on water droplet dispersion, evaporation and radiation blockage. The intent
770 of MaCFP is to expand the list of target experiments. It is also to keep re-visiting the initial
771 list for further insights into basic flow and combustion phenomena and for additional tests with
772 coarse-grained simulations.

773 In closing, the organizing committee of the MaCFP Working Group has now started prelimi-
774 nary discussions for the organization of a second workshop. Interested individuals/organizations
775 are encouraged to contact the committee [3] in order to join MaCFP, influence the selection of new
776 target experiments, participate in discussions on the structure, format and scope of MaCFP, partic-
777 ipate in discussions on the location and time of the second workshop. Note that the experimental
778 and computational databases corresponding to Cases 1-5 are hosted on the MaCFP repository [1]
779 and are available to the fire research community as reference data for future experimental and/or
780 computational studies. The MaCFP Working Group is committed to continuously update the
781 repository.

782 **3. Condensed Phase Subgroup**

783 *3.1. Objectives*

784 The production of combustible gases by burning materials is typically the rate-limiting process
785 in the growth of fire. A quantitative understanding of this process is therefore essential for advancing
786 our ability to predict and mitigate fire development. Unfortunately, measurement and modeling
787 efforts carried out in this field by various research groups tend to be poorly coordinated. Little
788 agreement exists as to what constitutes best practices and standards in data collection and model
789 development. The purpose of the condensed phase subgroup of the MaCFP Working Group is to
790 facilitate data and model sharing among researchers in order to improve predictions of thermal
791 decomposition and pyrolysis in fire. The work of this subgroup, in conjunction with the work of the
792 gas phase subgroup, is expected to lead to fundamental progress in fire modeling. It is envisioned
793 that the two subgroups will collaborate to make quantitative predictions of the combined gas-solid
794 phase processes that determine flame spread. It is also envisioned that the two subgroups will hold
795 joint workshops every two or three years.

796 The condensed phase subgroup shares the central objective of MaCFP “to target fundamental
797 progress in fire science and to advance predictive fire modeling.” The specific objectives of the
798 subgroup will focus on the development, calibration, verification, and validation of predictive models
799 of thermal decomposition and pyrolysis. To this end, the subgroup plans to:

- 800 • Develop several alternative formats for experimental data sets that carry sufficient information
801 to enable parameterization of pyrolysis models for a given material.
- 802 • Develop a set of requirements for data set quality and completeness and organize a committee
803 of experts that will review the submissions to the repository to ensure that they are compliant
804 with these requirements.
- 805 • Incorporate compliant data sets into the existing MaCFP data repository [1].
- 806 • Create a database of pyrolysis property sets that are generated from the experimental data
807 sets. Each pyrolysis property set will be required to be accompanied by a demonstration of
808 how well it captures the data on the basis of which it was calibrated and validated.
- 809 • Develop a set of minimum requirements for numerical pyrolysis simulation codes.
- 810 • Organize a discussion group focused on unresolved issues in pyrolysis modeling.

811 The scientific topics covered by the condensed phase subgroup will include:

- 812 • Kinetics and thermodynamics of the condensed phase decomposition reactions.
- 813 • Properties and composition of gaseous pyrolyzates.
- 814 • Heat and mass transfer in the condensed phase.
- 815 • Physics and chemistry of the gas-condensed phase interface including the topics of oxidative
816 pyrolysis and interactions with the surface flame.
- 817 • Coupled thermal and mechanical behavior of pyrolyzing solids including intumescence and
818 melt flow.

819 *3.2. Summary of the Planning Meeting*

820 As explained in section 1, in addition to being a first technical meeting for the gas phase
821 subgroup, the June-2017 MaCFP workshop served as a planning meeting for the condensed phase
822 subgroup. The planning meeting featured an introductory presentation by the co-chairs of the
823 condensed phase subgroup, followed by seven invited presentations and two periods for an open
824 discussion. Hard copies of the presentations can be found in [4].

825 The introductory presentation [69] presented an overview of the motivation, purpose, and goals
826 of the condensed phase subgroup. It was emphasized that fire phenomena can only be predicted
827 with robust coupling between condensed and gas phase models. Consequently, it will be necessary
828 for the two subgroups of MaCFP to work closely in the planning and analysis of validation data
829 as well as in subsequent model development. Several of the challenges associated with condensed
830 phase fire physics were mentioned. Overcoming these challenges requires systematic verification
831 and validation of condensed phase models. Several concepts from validation and verification were
832 reviewed including the so-called “validation pyramid” as a heuristic for systematically validating
833 complex models via sequential validation of various submodels. The International Workshop on
834 Measurement and Computation of Turbulent Nonpremixed Flames (known as the TNF workshop)
835 was mentioned as a model for organizing the condensed phase subgroup’s activities. The refer-
836 ence to the TNF model led to a brief description of a proposed plan for the subgroup’s work as
837 presented in the White Paper [3] prepared by the co-chairs prior to the workshop. The core of
838 the proposal is to facilitate communication between experimentalists and modelers by providing
839 web-based management of four elements: (1) experimental data; (2) numerical models; (3) param-
840 eter sets and associated comparisons between model predictions and experimental data; and (4)
841 a discussion forum. For each of these elements, the presentation provided a brief explanation as
842 well as some proposed constraints. First, the experimental data should initially focus on scenarios
843 in which flaming is not present so that condensed phase physics may be isolated. This data will
844 need to follow some requirements for formatting and review. Second, the numerical models should
845 be open source, well-documented, and include at least heat transfer and decomposition reaction
846 kinetics. Third, the parameter sets and comparisons should be complete and the link between the
847 underlying data and the parameter values should be specified. The Fire Dynamics Simulator (FDS)
848 Validation Guide [29] was mentioned as a good example of such comparisons. Finally, several topics
849 such as missing experimental data, needed model developments, and computational challenges were
850 suggested for the discussion forum. It was noted that a successful discussion forum will require

Table 1: Material properties required by state-of-the-art computational pyrolysis models.

Kinetic	Thermodynamic	Transport
Pre-exponential factors	Specific heat capacities	Thermal conductivities
Activation energies	Heats of decomposition reactions	Emissivities
Stoichiometric coefficients	Heats of combustion	Absorption coefficients
		Mass diffusivities

851 sustained community participation.

852 The first invited presentation [70] began with an overview of different condensed phase models,
 853 from early heat-transfer-based analytical models for ignition up to modern computational pyrolysis
 854 solvers such as FDS [31], Gpyro [71] and ThermaKin [72]. These modern computational models
 855 rely on a relatively large number of material properties used to characterize pyrolysis behavior. A
 856 list of common material properties used in computational pyrolysis models is provided in Table 1.
 857 Identifying values for these many parameters presents a challenge especially as the values can change
 858 significantly as a material heats and decomposes. The remainder of the presentation focused on
 859 describing a procedure for determining the kinetic and thermodynamic properties of materials
 860 developed at the University of Maryland [73]. This procedure relies on data from three milligram-
 861 scale experiments: (1) thermogravimetric analysis (TGA) for decomposition reaction kinetics; (2)
 862 differential scanning calorimetry (DSC) for heat capacities and heats of decomposition reactions;
 863 and (3) microscale combustion calorimetry (MCC) for heats of combustion of gaseous pyrolyzates.
 864 The presentation described these experiments and procedures for extracting material properties
 865 from the appropriate data. Throughout the discussion, data for poly(butylene terephthalate) (PBT)
 866 was used as an example.

867 The second invited presentation [74] discussed current work on validating models of solid react-
 868 ing materials. Surface temperature measurements using thermophosphors are being explored, and
 869 datasets for solid reactive materials are being generated. The focus at Sandia National Laboratories
 870 (Sandia) is on the high heat flux regime. A number of test facilities are available for high heat flux
 871 ignition experiments including the Sandia solar furnace which can provide a heat flux of 5 MW/m².
 872 In addition to experimental work, Sandia is developing a code for fire modeling (Fuego [22]) and a
 873 code for reacting solid materials (Aria).

874 The relationship between gas and condensed phase physics was discussed in the third invited
 875 presentation [75]. The large number of physical processes occurring at the solid-gas interface were

876 enumerated. It was emphasized that ignition and flame spread are significantly influenced by
877 the details of transport and chemistry occurring at the interface. A review of boundary layer
878 theory was provided as well as a discussion of the reacting boundary layer theory of Emmons.
879 The presentation concluded by highlighting the need to develop better models of turbulence, heat
880 transfer, and combustion in the near-wall region of the boundary layer to account for chemical and
881 blowing effects corresponding to pyrolysis.

882 The fourth invited presentation [76] provided a discussion of the solid model implemented in
883 FDS [18]. The presentation began by presenting a schematic of the physical processes involved in
884 burning materials and emphasized the multi-scale nature of the problem. The governing equations
885 for condensed phase species and energy conservation as well as the pore gas conservation equa-
886 tions for species, energy and momentum, were presented. The need to limit the model to include
887 only the important physics was noted. Following on from this point, the presentation listed the
888 major assumptions made by the FDS solid model. Specifically, the FDS solid model assumes one-
889 dimensional transport, no mass accumulation (and therefore instantaneous mass transfer), thermal
890 equilibrium between gases and solids, and assumes that a heat of reaction may be used to account
891 for the energy contribution of the decomposition reactions. The FDS solid model is regularly ver-
892 ified (in terms of heat conduction, radiation, mass conservation, and reaction rate) and validated
893 (in terms of mass loss rate and heat release rate for burning polymer slabs). Several special topics
894 for future pyrolysis model development were mentioned including spectral radiation, shrinking and
895 swelling, pressure build-up, multi-dimensional effects, and solid mechanical considerations associ-
896 ated with fracturing of char. The presentation concluded by suggesting that these special topics
897 might be appropriate for further exploration by the MaCFP condensed phase subgroup.

898 The challenge of coupling condensed and solid phase models was explored in the fifth invited
899 presentation [77]. In contrast to the multi-experiment approach developed at the University of
900 Maryland [70], FM Global calibrates all material properties using data from a single experiment,
901 namely the fire propagation apparatus (FPA). A one-dimensional model with a single-step Arrhe-
902 nius reaction is then fit to the FPA data using optimization with the shuffled complex evolution
903 (SCE) algorithm. The resultant pyrolysis properties are then used as inputs in a CFD fire model
904 (FireFOAM [20]) to simulate full-scale fire scenarios. Validation of this approach has been per-
905 formed for several additional FPA scenarios. An important application of fire models for FM
906 Global is understanding fire spread in warehouse rack storage. FireFOAM has been used to predict
907 heat release rate in 3-tier, 5-tier, and 7-tier rack storage of cardboard boxes using properties ob-

908 tained by the FPA/SCE material property calibration procedure. The presentation also discussed
909 applications involving boxes of plastic cups and large rolls of paper. The paper rolls present a
910 unique challenge in that delamination of outer layers of paper had to be accounted for. Several
911 lessons were provided in conclusion. First, coupling of gas and condensed phase models is neces-
912 sary for real-world problems, and the appropriate level of model complexity is determined by the
913 problem. Second, validation needs to occur both for the decoupled and coupled models at multiple
914 scales.

915 The sixth invited presentation [78] discussed recent work on assessing the appropriate level of
916 model complexity. Beginning with a clear statement of the goal to “up-scale” from fundamental
917 physics and chemistry to real fire behavior, the presentation laid out some of the many challenges
918 in the path of achieving that goal. One of those challenges is choosing the appropriate level of
919 model complexity. The number of parameters in pyrolysis models can vary from just a few to
920 over 30 for some of the more detailed models in existence. The problem of complexity is one of
921 finding the minimum number of parameters required to attain an acceptable level of error. In
922 the recent work presented in Refs. [79, 80], this problem has been addressed by systematically
923 decreasing model complexity used to predict the pyrolysis of a vertical slab of PMMA exposed to
924 varying levels of heat flux and oxygen. It was found that it is not helpful to increase the complexity
925 of the chemical model unless a sufficiently complex model of heat transfer is used. Furthermore,
926 additional complexity corresponds to increased uncertainty, and so complex models should only be
927 used in the presence of sufficient, high quality data.

928 Finally, the seventh invited presentation [81] provided an overview of pyrolysis modeling with
929 Gpyro. Gpyro is an open-source three-dimensional pyrolysis model with user-specified complexity.
930 Additionally, Gpyro may be coupled to FDS with some limitations (*e.g.* Cartesian geometries,
931 no shrinkage or swelling, and no burn-away). Current work involves coupling to ABAQUS for
932 mechanical calculations. A critical part of any pyrolysis solver is the material property models
933 allowed. Gpyro treats material properties as weighted sums of species properties with a power
934 law dependence on temperature. Additionally, permeability and thermal conductivity may be
935 anisotropic which can be important for materials such as wood. After going through the form
936 of the conservation equations, the presentation provided some details on the numerical schemes
937 employed by Gpyro. The time-stepping is fully implicit to ensure stability of the solution, and
938 an alternating direction tri-diagonal matrix algorithm is utilized for speed. Several verification
939 cases have been carried out including for a sphere with internal heat generation. As an example of

940 the power of detailed pyrolysis modeling, the presentation gave the example of wood pyrolysis at
941 both small and large external heat fluxes, with the “fast” case producing significantly more tar as
942 compared to the “slow” case.

943 The invited presentations served to establish a foundation for the two periods of open discussion
944 that were scheduled in the workshop. These open discussion periods were crucial for fielding input
945 from the research community at large. The issue of heating rate in small scale experiments was
946 discussed. Some believe that the heating rates used in small-scale tests should emulate realistic fire
947 heating rates, but it was noted that chemical reaction kinetics generally depend on temperature
948 (not heating rate) and increasing the heating rate does not significantly change the temperature
949 range over which solid decomposition reactions take place. Furthermore, high heating rates can
950 lead to temperature and species concentration gradients which prevent meaningful interpretation of
951 results of small-scale tests such as TGA. Several participants suggested that TGA should be coupled
952 with gas analysis (e.g., FTIR) as this information could be important for the gas phase physics.
953 Similarly, the impact of oxygen concentration should be explored further in order to model the
954 transitional regimes of ignition, spread, and extinction in contrast to steady burning. For all small-
955 scale tests, it was suggested that it should be necessary to precisely describe the range of validity
956 of the parameters as a consequence of how they are determined. Much of the open discussion time
957 was devoted to identifying appropriate validation data sets. Some participants suggested that it is
958 important to have large-scale data (such as the FM Global parallel panel test or the standard room
959 corner tests) early on in order to better guide subsequent model development and experimentation.
960 This would not negate the necessity of small-scale tests for model parameterization or validation
961 of sub-models, but inversely, this illustrates the pertinence of the up-scaling approach. Another
962 issue that arose is the appropriate selection of test materials. A balance should be struck between
963 simplicity for modeling purposes and real-world application.

964 *3.3. Future Plans*

965 Future steps include the development of a digital archive dedicated to the condensed phase
966 subgroup, possibly using the same platform as the gas phase subgroup [1]. In parallel, the standards
967 for the experimental data sets will be established. The following standards are proposed:

- 968 • Each studied material must have clearly defined chemical composition. The material’s phys-
969 ical attributes, such as color, initial density and thickness, geometry of reinforcement (in the

970 case of structural composites), must be provided. The material should be readily available,
971 preferably, from multiple distributors.

972 • The material should be conditioned prior to all experiments in a well-defined atmosphere with
973 these conditions specified. For hydrophilic materials, the initial moisture content should be
974 reported.

975 • The experiments used to determine properties may consist of milligram-scale and/or gram-
976 scale tests. Milligram-scale tests (such as TGA) are expected to be conducted under thermally
977 thin conditions, *i.e.* conditions for which the sample temperature is spatially uniform and
978 resolved in time. Gram-scale tests (such as FPA gasification experiments [82]) are expected
979 to be conducted under non-thermally thin conditions, *i.e.* conditions for which transport
980 properties have significant impact on the measured quantities. Gram-scale tests must have
981 well-defined thermal boundary conditions, more specifically, heat fluxes incident on all sample
982 surfaces should be specified as a function of surface temperature. If the material sample is
983 mounted onto thermal insulation, the properties of this insulation must be provided. The
984 composition of the gaseous environment inside all test apparatus must be defined. For all
985 tests, the size and mass of the samples must be fully specified.

986 • Each data set may contain either milligram- and gram-scale test results or, alternatively, only
987 gram-scale test results. In the latter case, the results from multiple experiments performed at
988 a range of heating conditions must be reported and include time-resolved sample mass as well
989 as sample temperature measurements (at the surface and/or at an in-depth location). The
990 conditions of the tests (for example heating rate, temperature range, percentage of oxygen in
991 the atmosphere) must be defined.

992 • Heat of combustions of gaseous pyrolyzate produced by the material must be measured using
993 a Cone Calorimeter, FPA, or Microscale Combustion Calorimeter.

994 • Additional data, including chemical composition of the gaseous pyrolyzate, and thermal con-
995 ductivity, emissivity, radiation absorption coefficient and mass diffusivity of the solid, are
996 desirable but not required.

997 • All experimental data must contain information on their uncertainties.

998 Multiple experimental data sets for the same material will be allowed into the repository, pro-
999 vided that each of them satisfies all established requirements. One key requirement for each py-
1000 rolysis property set, which will be generated from the experimental datasets, is a demonstration of
1001 how these property values capture all data in at least one experimental dataset, *e.g.* if the dataset
1002 contains the results of controlled-atmosphere cone calorimetry experiments and TGA, the developer
1003 of the pyrolysis property set will be required to produce predictions of both of these experiments
1004 and provide input files that were used to generate these predictions. Quantitative criteria will be
1005 developed to characterize the quality of each prediction.

1006 It is proposed that, initially, experimental data sets will be developed for relatively simple
1007 materials that are isotropic in nature and do not exhibit complex mechanical behavior such as
1008 melt flow, delamination or intumescence. Examples of such materials include cast poly(methyl
1009 methacrylate) and high-impact polystyrene. Demonstrating that the pyrolysis property sets can be
1010 used to successfully predict compartment-scale fire growth will be a longer term goal of this effort.
1011 One potential target geometry for the full-scale experiments is upward and lateral flame spread in
1012 a flammable corner, which is realized in several flammability standards [83–85]. It is proposed that
1013 well-instrumented versions of these standard experiments be carried out to serve as a modeling
1014 target for comprehensive gas and condensed phase models of fire growth.

1015 In closing, the co-chairs of the condensed phase subgroup of the MaCFP Working Group
1016 have now started discussions for the organization of a second workshop. Interested individu-
1017 als/organizations are encouraged to contact the co-chairs [3] in order to participate in preparations
1018 for the workshop and in the construction of the digital archive described above.

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