

Effect of ultrasonic homogenization on the Vis/NIR bulk optical properties of milk

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Abstract

The size of colloidal particles in food products has a considerable impact on the product's physicochemical, functional and sensory characteristics. Measurement techniques to monitor the size of suspended particles could, therefore, help to further reduce the variability in production processes and promote the development of new food products with improved properties. Visible and near-infrared (Vis/NIR) spectroscopy is already widely used to measure the composition of agricultural and food products. However, this technology can also be consulted to acquire microstructure-related scattering properties of food products. In this study, the effect of the fat globule size on the Vis/NIR bulk scattering properties of milk was investigated. Variability in fat globule size distribution was created using ultrasonic homogenization of raw milk. Reduction of the fat globule size resulted in a higher wavelength-dependency of both the Vis/NIR bulk scattering coefficient and the scattering anisotropy factor. Moreover, the anisotropy factor and the bulk scattering coefficients for wavelengths above 600 nm were reduced and were dominated by Rayleigh scattering. Additionally, the bulk scattering properties could be well ($R^2 \geq 0.990$) estimated from measured particle size distributions by consulting an algorithm based on the Mie solution. Future research could aim at the inversion of this model to estimate the particle size distributions from Vis/NIR spectroscopic measurements.

Keywords

Food emulsion, Milk, Ultrasonic homogenization, Vis/NIR spectroscopy, Light scattering, fat globule size distribution

1. Introduction

Many products in the food industry, including milk, ice cream, mayonnaise, jam, etc., are colloids or have been produced via colloids. Dispersed particles determine, through their concentration, distribution, size and structure, the physicochemical and organoleptic quality properties of the product. Additionally, the complex interactions among colloidal particles and with other ingredients will contribute to the rheology, texture, stability, appearance, and many other food characteristics [1]. Techniques to measure the particle properties of the colloid systems allow for better understanding these complex relations and further improve food processing techniques and recipes. Furthermore, implementation of such sensing-technology in a food plant could enable online monitoring of the colloids or the derived products during the production process and promote early detection of an altering product quality [2].

Milk is a natural colloidal emulsion of fat globules, as well as a hydrocolloid suspension of casein micelles, both dispersed in a water-based solution [3]. The fat globules, with a size ranging from 0.1 to 15 μm , result from the assembly of small intracellular lipid droplets, followed by an apocrine secretion through the apical membrane of the epithelial cells in the mammary gland [4]. The volume fraction of these fat globules and casein micelles, the particle size and shape and particle-particle interactions contribute significantly to the quality

50 properties of dairy products [5]. Among these influential factors, the milk fat globule (MFG)
51 size was found to have an important impact on the physicochemical, functional and sensory
52 characteristics of the milk and derived products [5–9]. The size of the native fat globules is
53 correlated with the fatty acid composition and bioavailability of the lipids [10–18]. Since the
54 size of the globules, and more specifically the amount of fat globule membrane, has an impact
55 on the enzymatic activity in the milk, it will affect the colloid stability. Furthermore, large fat
56 globules cream faster and make skimming easier [13,19]. Centrifugation, gravity separation
57 and microfiltration are methods to separate native MFGs of different sizes to create products
58 with new, better and/or standardized properties [13,19–22]. High-pressure or ultrasonic
59 homogenization can also reduce the size of the fat globules by disrupting them into smaller
60 ones [23–28]. However, as the MFG membranes are disrupted too and casein micelles and
61 whey proteins adsorb at the increased surface area, the properties for homogenized MFGs are
62 different compared to small native MFGs [10,11,7,20]. Because of the impact of these
63 techniques on the size of the fat globules and accordingly the food colloid properties, it is
64 important to monitor these processes to ensure a high quality product with minimal intra- and
65 inter-batch variability [6,29]. This is also the case for raw milk as there is a large variability in
66 MFG quantity and size distribution, varying between milk of different cows, within individual
67 cows between different milkings and quarters, and within individual quarter milkings between
68 different times of sampling [6,29–31]. Hence, online prediction of the fat content and MFG
69 size distribution, during the milking process, would allow for the separation of raw milk with
70 different colloidal properties, depending on the needs of the food industry [32,33]. This could
71 reduce the number of process steps in a food plant, promote the development of new food
72 products with extra functional properties and upgrade the value of raw milk [29].
73 Additionally, as the size of MFG in milk from infected mammary glands is significantly
74 increased, MFG size information obtained on the dairy farm could give insight into the udder
75 health status of each individual cow and udder quarter [34,35].

76 Optical measurement techniques are frequently used to monitor processes in food
77 industry and agriculture. Among them, spectroscopy studies the interaction between
78 electromagnetic radiation (EMR) and the product of interest. Because of the relatively deep
79 penetration, relaxing the requirements for sample preparation, visible (Vis) and near-infrared
80 (NIR) spectroscopy is often preferred. It combines a fast and non-destructive measurement
81 principle with robust and cost-efficient sensor technology available nowadays. Originally,
82 Vis/NIR spectroscopy was used to study the wavelength-dependent absorption by molecules
83 in the product. The measured absorption relates to the number of molecules that absorb at that
84 wavelength and the sample thickness. Therefore, it can be consulted to predict the product's
85 composition through the Beer-Lambert law. As most food and agricultural products contain
86 local non-uniformities, the radiation is forced to deviate from its straight trajectory and the
87 travelling path is increased to an unknown extent. These non-uniformities can be polymer,
88 metal or mineral particles in industrial systems; fat globules in emulsions; cells, cell
89 organelles, air pores and fibrous structures in biological tissues; or pores, air bubbles and
90 crystals in processed foods, etc. These elastic scattering processes complicate the prediction of
91 the composition from measured spectra as the measurements have to be corrected for this non-
92 linear interference [32]. Moreover, samples with scattering properties that are out of the range
93 compared to the samples used in the calibration will result in inferior predictions [36]. On the
94 other hand, the scattering properties could also be exploited as a valuable source of
95 information because the overall scattering level and the angular- and wavelength-dependency
96 are determined by the physical microstructure of the product. For emulsions, the scattering
97 properties are related to the quantity, size distribution and material properties (refractive
98 index) of the suspended spherical particles [6,23–26,37].

99 In the last decade, several models and measurement techniques based on Vis/NIR
100 spectroscopy have been proposed to quantify the absorption and scattering properties for
101 homogeneous food products [38–42]. The derived bulk absorption coefficient μ_a (cm^{-1}) is
102 defined as the probability of absorption per unit infinitesimal path length and can be related to
103 the product composition without the need for (empirical) scatter corrections. The scattering,

104 on the other hand, can be described with the bulk scattering coefficient μ_s (cm⁻¹) and the
105 angular scattering pattern: the normalized scattering phase function $p(\theta)$. μ_s describes the
106 probability of scattering per unit infinitesimal path length, analogous to μ_a , while $p(\theta)$ gives a
107 detailed description of the normalized scattering probability as a function of the scattering
108 angle θ . As $p(\theta)$ is often too complex to represent and interpret, only the first moment of $p(\theta)$,
109 defined as the scattering anisotropy factor g , will generally be considered and represents the
110 mean cosine of the scattering angle. The anisotropy factor for biological tissues and emulsions
111 in the Vis/NIR range varies between isotropic Rayleigh scattering ($g = 0$) and complete
112 forward scattering ($g = 1$) [43]. Extraction of these bulk scattering properties from Vis/NIR
113 spectroscopic measurements, and prediction of the particle size distribution (PSD) of the
114 measured emulsion from these properties could create a whole new field of applications for
115 Vis/NIR spectroscopy. As Vis/NIR spectrometers are already well established for
116 compositional analysis in the agro-food industry, extending their functionalities with PSD
117 estimations would further promote the implementation of this technology [6,44].

118 In order to derive microstructure information (e.g. PSD) from measured bulk scattering
119 properties (μ_s and $p(\theta)$), the complex relation between those two has to be unravelled by
120 solving Maxwell's equations for electromagnetic theory [43]. Since the scattering particles in
121 an emulsion are almost perfectly spherical, the equations can be evaluated with the Mie
122 solution. The Mie solution for Maxwell's equations describes the electromagnetic field inside
123 and outside a spherical scatterer for given particle properties. Nowadays, algorithms that rely
124 on the Mie solution are available to simulate the bulk optical properties (μ_a , μ_s and $p(\theta)$ or g)
125 for polydisperse spherical particles in an absorbing host medium, taking into account the
126 particle's volume concentration and size distribution, the complex refractive indices of
127 particle and medium, and the considered radiation wavelength [37,45]. Consequently,
128 inversion of these algorithms would clear the way for PSD estimations from measured bulk
129 scattering properties [6].

130 However, these algorithms build on a number of assumptions which may not always be
131 valid for food colloid systems. Apart from the assumption of perfect sphericity of the
132 particles, these algorithms also assume that scattering events are independent. Under this
133 condition, μ_s follows a linear increase (through the origin) with increasing particle quantity,
134 while the normalized scattering phase function and the anisotropy factor are independent. In
135 practice, however, the independent scattering relation becomes invalid when the particle
136 volume concentration is above a certain threshold. It has been shown that this limit depends
137 on the radiation wavelength, refractive indices of the particles and host medium and the
138 particle size, shape and polydispersity [46,47]. Beyond this limit, scattering particles are close
139 enough such that individual scattering events can influence each other. This dependent
140 scattering typically reduces the bulk scattering coefficient and the anisotropy factor compared
141 to what could be expected from the independent scattering relations [48–50].

142 To evaluate the concept of PSD estimation from the bulk scattering properties of milk, it
143 is crucial to validate the accuracy and robustness of the algorithms that relate the bulk
144 scattering properties to the particle size distribution of milk. Accordingly, this is the main
145 objective of the present study. More specifically, the effect of the fat globule size and quantity
146 on the bulk optical properties (BOP) of milk was investigated by measuring milk samples
147 with the 'golden standard' method and comparing them to the BOP simulated from the
148 measured particle size distributions.

149 This manuscript is an extended and improved version of a conference proceeding
150 presenting the preliminary results [40]. Moreover, the PSD measurements were measured with
151 a laser diffraction apparatus with improved sensitivity for submicron particles (paragraph 2.2),
152 while the accuracy of the unscattered transmittance measurements was further enhanced
153 (paragraph 2.3).

154 **2. Materials and methods**

155 *2.1 Milk samples and ultrasonic homogenization of milk*

156 A fresh, raw bulk milk sample was collected from a Belgian dairy farm. The major
157 components were measured with a Milkoscan FT+ (Foss A/S, Hillerød, Denmark) milk
158 analyzer according to ISO 9622:2000. The fat, crude protein, lactose and water content were
159 respectively 4.12%, 3.53%, 4.55% and 89.7% (all *w/v*). To create milk samples with different
160 PSDs, the original raw milk sample was heated up to $37 \pm 1^\circ\text{C}$ in a water bath, gently shaken
161 and split into 7 subsamples of 20 ml (polypropylene cup, 28 mm diameter and 68 mm high) to
162 be homogenized with ultrasound for respectively 0 (raw), 20, 60, 120, 300, 600 and 1200 s.
163 Ultrasound waves were generated by a Vibra-Cell VCX400 (400 Watt, 20kHz) equipped with
164 a CV26 converter, a standard 13 mm horn and a tapered 3 mm microtip (all Sonics &
165 Materials Inc., Danbury, USA). The microtip was immersed into the milk at a depth of 5 ± 1
166 mm. The power level of the ultrasound generator was fixed at 20% (80 Watt) to prevent the
167 milk from foaming, but still maintaining a high homogenization efficiency. During
168 homogenization, the cup with the milk sample was cooled in a water bath at $\pm 24^\circ\text{C}$ to
169 maintain a constant sample temperature of $37 \pm 1^\circ\text{C}$. Consequently, denaturation of the
170 proteins was prevented and a high homogenization efficiency was preserved [27]. For
171 additional comparison, a commercial half-skimmed, homogenized (further referred to as
172 'Skim') milk sample (Joyvalle, Belgium), with a fat and crude protein content of respectively
173 1.5% and 3.5% (both *w/w*), was also considered in this study.

174 The PSDs and bulk optical properties (BOP) were measured for all 8 samples (next
175 paragraphs). To test the validity of the independent scattering condition, the BOP of the raw
176 milk sample and the milk samples homogenized for 120 and 1200 s were, next to their pure
177 state ($1/1$), also measured after two- ($1/2$) and three-fold dilution ($1/3$) with deionized water.

178 *2.2 Milk particle size distribution measurements with laser diffraction*

179 The size distribution of fat globules in raw and homogenized milk is generally measured
180 with laser diffraction (LD) after dissociating the suspended casein micelles [10,11,14–
181 16,18,7–9,20,28,29,51–62]. Nevertheless, studies showed that the presence of casein micelles
182 does not interfere with the size distribution of the fat globule fraction as determined with LD
183 [28,59,62]. Moreover, the protein dissociating additives might also affect the actual fat
184 globules as they can disrupt the MFG membrane proteins [59]. Therefore, the fat globule size
185 distribution can be determined through LD without the need for casein micelle dissociation
186 [13,6,63,64]. This is especially true for raw milk as there is (nearly) no overlap in the fat
187 globule and casein micelle size.

188 The size distribution of milk casein micelles is usually determined with dynamic light
189 scattering (DLS) after removing the fat globules through skimming [65–74]. However, the
190 particle size measured with DLS relates to the hydrodynamic radius of the dynamic
191 hydrated/solvated particle. The latter depends, next to the actual particle radius, also on the
192 interaction between the particles and the solvent. As the surface of the casein micelles is
193 covered with highly hydrophilic layer of κ -casein, the hydrodynamic radius is significantly
194 larger than the particle-core radius [67]. In this study, however, the particle-core size is of
195 interest as this one relates to the static light scattering and consequently the BOP of the
196 suspension. In contrast, the PSD measured with LD techniques is related to the actual particle-
197 core size. Further on, the newest generation LD devices are sufficiently accurate to also
198 measure the size of sub-micron colloidal particles as small as 10 nm [75,76]. As a result, LD
199 is used in numerous recent studies to determine the size of casein micelles in skimmed milk
200 [56,60,77–85]. In these studies, the fat globules were removed from the milk through
201 centrifugation or microfiltration in order to measure the PSD of only the casein micelle
202 fraction. Nevertheless, the measured PSD of the casein micelle fraction, as determined with
203 LD techniques, was found to be nearly independent of the presence of fat globules
204 [56,59,60,84]. Additionally, the centrifugation and/or microfiltration process step might affect

205 the actual casein micelles size distribution. Consequently, it was concluded that the casein
206 micelle size distribution can be determined through LD without removal of the fat globules,
207 especially in raw milk [56,59,60,84].

208 In this study, the size distribution of the fat globules and casein micelles was measured
209 simultaneously with LD. The Hydro EV sample dispersion unit of the LD equipment
210 (Mastersizer 3000, Malvern, UK) was filled with 600 ml of deionized water and the stirring
211 speed was set at 2400 rpm. Milk was added drop-wise until the red laser (633 nm) obscuration
212 due to light scattering was $7\pm 1\%$ (ISO 13320:2009). Measurements were obtained directly
213 after reaching the intended obscuration level as casein micelles start to break down due to low
214 calcium concentrations at high dilution ratios [86]. Each sample was loaded two times to the
215 LD instrument and 5 measurements were obtained per loading. The PSD of the fat globules
216 and casein micelles was estimated from the measured angular scattering intensities by fitting
217 the Mie model for spherical particles (Mastersizer 3000 software, Malvern, UK). The Mie
218 model uses the refractive indices for the milk particles and milk serum at the wavelengths of
219 the instrument's light sources, respectively 633 (red) and 470 nm (blue) [57,87]. Next, the
220 goodness of fit was evaluated based on the residual and the weighted residual. According to
221 ISO 13320:2009, a weighted residual above 1% might indicate an erroneous measurement
222 and/or the use of incorrect refractive indices for the sample and dispersant. In this study, only
223 PSD's were considered if the residual and the weighted residual of the fit were below
224 respectively 0.8% and 0.6%. Next, the repeatability of the measurement and fitting procedure
225 was evaluated based on the coefficient of variation (CV) for the Dv10, Dv50 and Dv90, being
226 the particle diameter for which respectively 10, 50 and 90% (v/v) of the particles is smaller
227 than or equal to (ISO 13320:2009). Only if the CVs for the Dv10, Dv50 and Dv90 were below
228 2% for all 10 (2 x 5) measurements on a sample, the average PSD was calculated and
229 considered in the further analysis.

230 The PSD of casein micelle fraction and fat globule fraction could clearly be separated for
231 raw milk, but were overlapping if fat globules were disrupted due to homogenization. The
232 PSD of the casein micelles itself could be determined from the first mode ($< 0.8 \mu\text{m}$ diameter)
233 of the bimodal PSD of the raw milk sample. It was assumed that the casein micelle volume
234 fraction and size distribution were not influenced by the ultrasonic homogenization process.
235 Accordingly, the fat globules size distributions for all 8 milk samples could be obtained from
236 the originally measured PSDs after subtracting the casein micelle size distribution.

237 Additionally, the fat globules for the raw milk sample and the milk samples homogenized
238 for 120 and 1200 s were also inspected with a Leitz Diaplan microscope (Leica Microsystems
239 GmbH, Wetzlar, Germany) at 1250x magnification.

240 *2.3 Measurement of bulk optical properties for raw and homogenized milk samples*

241 Double integrating sphere (DIS) and unscattered transmittance measurements are
242 considered to be the 'golden standard' to determine the BOP (μ_a , μ_s and g) for thin samples of
243 turbid media and were therefore used in this study [88]. The diffuse reflectance and total
244 transmittance were measured simultaneously on the milk sample loaded in a borosilicate
245 cuvette with a path length of 600 μm and positioned between the two integrating spheres, both
246 equipped with a Vis and NIR detector. The unscattered transmittance was measured in a
247 separate path with the Vis and NIR detector positioned 1.5 m behind the sample, which was
248 loaded in a similar cuvette with a path length of 170 μm . A series of slits between sample and
249 detector further reduced the number of scattered photons captured in the unscattered
250 transmittance signal [48,88]. The sample illumination was especially designed to obtain high
251 signal-to-noise spectra in the 500 – 2250 nm wavelength range for very turbid media. It
252 consists of a supercontinuum laser light source (500 – 2250 nm, 4 Watt optical power) in
253 combination with a high-precision monochromator. For more details on the measurement
254 principle, the setup, its validation and the calibration and measurement procedure, the reader
255 is referred to [88]. All sample spectra were measured from 500 until 1900 nm in steps of 10
256 nm by automated scanning of the monochromator. The samples were thoroughly stirred
257 before they were measured at room temperature.

258 The inverse adding doubling (IAD) routine developed and optimized by Prahl [89] was
259 consulted to obtain the BOP values from the diffuse reflectance and total and unscattered
260 transmittance measurements. Apart from the three measurements, also the real refractive
261 index of the sample had to be provided to the algorithm and was calculated with the equation
262 proposed in [87]. Because of significant contribution of scattered photons to the unscattered
263 transmittance signals, no BOP estimation could be established for raw milk at wavelengths
264 below 600 nm and for milk samples homogenized for 20, 60, 120 and 300 s at wavelengths
265 below 550 nm [48,88].
266

267 *2.4 Simulation of the bulk optical properties for raw and homogenized milk samples*

268 The BOP of the commercial, raw and homogenized milk samples were also calculated
269 from their respective PSD with an algorithm to simulate the BOP of polydisperse suspensions
270 of spherical particles in an absorbing host medium [45]. In this algorithm, the Mie solution for
271 Maxwell's equations has been generalized to account for the effect of an absorbing host
272 medium on the particle's scattering properties. Polydispersity of the particle system is
273 supported by fractionalization of the provided PSD, while the number of intervals is
274 automatically optimized in an efficient iterative procedure. This procedure is straightforward
275 as it does not involve an inversion or fitting procedure. For more details on the algorithm and
276 its validation, the reader is referred to [45]. In this study, the algorithm was used to simulate
277 the BOP of the samples in the 500 – 1900 nm wavelength range, starting from the measured
278 PSDs and refractive indices obtained from literature.

279 The BOP for the fat globule fraction and casein micelle fraction were simulated
280 separately as their refractive indices differ significantly. Moreover, the fat globule size
281 distribution, fat volume concentration and fat and milk serum refractive indices were provided
282 as inputs for the algorithm to simulate the BOP for the fat globule fraction, while a similar
283 approach was followed for the casein fraction. The fat volume concentration was found to be
284 4.52% v/v, as derived from the fat content (paragraph 2.1) and the fat and milk density of
285 respectively 0.912 and 1.03 g/ml [87]. The casein volume concentration was 2.00% v/v,
286 calculated from the crude protein content (paragraph 2.1), by taking into account the average
287 casein fraction in crude protein (75.5% w/w) and the density of casein (1.33 g/ml) and milk
288 [87]. The real part of the refractive index spectra for milk fat, casein and milk serum was
289 obtained by adding a baseline to the Vis/NIR refractive index spectrum of water [90]. The
290 baseline was determined at 589 nm, a wavelength for which the refractive indices are known
291 for water, milk fat, casein and milk serum [87,90]. The imaginary part of the refractive index
292 for milk fat, casein and milk serum was calculated from the bulk absorption coefficients of
293 respectively fat, protein and a mixture of water, lactose and serum proteins [37,90–93]. The
294 calculated BOP for the fraction of fat globules and casein micelles were afterwards
295 recombined, similar to the recombination of different PSD fractions, to obtain the BOP of the
296 milk samples [45].

297 **3. Results and discussion**

298 *3.1 Particle size distribution of raw and homogenized milk samples*

299 PSDs were considered after they passed the quality inspections given in paragraph 2.2.
300 The final cumulative volume-frequency PSD for each of the 8 milk samples was calculated as
301 the average of the passed replicates and is presented for each sample in Fig. 1. For raw milk,
302 an explicit bimodal PSD can be noticed. Casein micelles with a size between 0.01 and 0.8 μm
303 and fat globules with a size between 0.8 and 12 μm represent respectively the first and second
304 mode. As the homogenization time increases, fat globule size is reduced, resulting in an
305 overlap with the casein micelles. For the commercial milk sample and for samples
306 homogenized over a period of 300 s and longer, no distinction in size can be made between
307 the casein micelles and fat globules.

308 The casein micelle size distribution, derived from the first mode of the bimodal PSD for
309 raw milk, was subtracted from all the measured PSDs, resulting in the cumulative volume-
310 frequency size distribution for the fat globules [Fig. 2]. This subtraction method resulted in
311 smooth monomodal PSDs, which indicates that the ultrasonic homogenization process had no
312 noticeable effect on the casein micelle volume fraction and size distribution [94], as was
313 assumed earlier (paragraph 2.2). Fig. 2 clearly demonstrates that the size of the fat globules
314 decreases with increasing homogenization time, which implicates that their quantity increases.
315 The PSD derived for the casein fraction is also presented in Fig. 2.

316 The most important properties describing the PSDs in Fig. 1 and Fig. 2 are listed in Table
317 1. Next to the percentiles Dv10, Dv50 and Dv90, also the surface area moment mean (D32)
318 and the volume moment mean diameter of the particles (D43) are presented. Before
319 subtraction of the casein micelle PSD [Table 1(a)], homogenization had nearly no effect on
320 the Dv10 as the 10% (v/v) smallest particles were mainly casein micelles. All other PSD
321 properties reduced steadily with increasing homogenization time. The half-skimmed (Skim)
322 milk sample contained fat globules of similar size compared to the milk sample homogenized
323 for 1200 s. However, the PSD width for the '1200 s' sample is wider [Table 1(b) and Fig. 2].

324 In respectively Fig. 3, Fig. 4 and Fig. 5, the microscopic images (1250x magnification) of
325 raw milk and milk homogenized for 120 and 1200 s are illustrated. They clearly demonstrate
326 the effect of ultrasonic homogenization on the size of the fat globules. After 1200 s of
327 ultrasonic homogenization, the fat globules are too small to be clearly noticeable with an
328 optical microscope.

329 3.2 Measurement of the bulk optical properties for raw and homogenized milk samples

330 The BOP for the 8 undiluted milk samples were obtained through IAD from DIS and
331 unscattered transmittance measurements. For all the raw and homogenized milk samples, the
332 measured bulk absorption coefficient spectra were very close ($R^2 \geq 0.968$) to each other and to
333 the μ_a spectrum of the water fraction (89.7% v/v) [90]. This is because the composition of
334 these samples is exactly the same, and water is by far the most important Vis/NIR absorbing
335 milk component [Fig. 6]. In respectively Fig. 7 and Fig. 8, the measured bulk scattering
336 coefficient spectra (μ_s) and scattering anisotropy spectra (g) are illustrated for the undiluted
337 milk samples. Both the μ_s spectra [Fig. 7] and the g spectra [Fig. 8] are highly susceptible to
338 the size (distribution) of the scattering fat particles. Ultrasonic homogenization for less than
339 one minute resulted in an increase of the μ_s for radiation wavelengths below 600 nm and a
340 decrease of the μ_s at longer wavelengths, while longer homogenization induced an overall
341 decrease of the sample's Vis/NIR μ_s spectrum. Furthermore, with reducing particle size, the μ_s
342 spectra became more dependent on the radiation wavelength (λ), approaching Rayleigh
343 scattering ($\sim \lambda^{-4}$) [43]. The μ_s spectrum for the half-skimmed milk sample was lower compared
344 to the '1200 s' sample, despite the similar fat globule size. This can mainly be explained by
345 the lower fat volume concentration and consequently a lower quantity of scattering fat
346 globules [23–26].

347 The anisotropy factor for the raw milk sample in the Vis/NIR range was relatively high
348 (0.934 – 0.960) and nearly independent of the radiation wavelength. With reducing fat globule
349 size, the anisotropy factor decreased and became more wavelength-dependent. This reduction
350 in g indicates that the light scattering was more isotropic, which is typical for smaller
351 scattering particles [43]. No accurate values for the anisotropy factor could be established for
352 the samples 'Skim' and '1200 s' for radiation wavelengths above 1000 and 1450 nm,
353 respectively. This was caused by diffuse reflectance values close to the noise level due to the
354 very high absorption by water ($\geq 89.7\%$ v/v) in combination with low scattering ($\mu_s < 10 \text{ cm}^{-1}$).
355 If scattering processes are independent of each other, the g spectrum should be independent
356 of the particle quantity [48]. However, despite their similar fat globule size and similar casein
357 micelle size [Table 1(b)], the g spectrum for the half-skimmed milk sample was lower
358 compared to the '1200 s' sample. Because the fat globule quantity was lower in sample
359 'Skim', while the casein micelle quantity was similar, the micelles would have contributed
360 more to the final g spectrum of sample 'Skim' compared to '1200 s'. As the casein micelles

361 were smaller than the fat globules [Table 1(b) and Fig. 2], this resulted in a lower g spectrum
362 and more isotropic scattering [Fig. 11]. Consequently, the g spectrum for sample ‘Skim’ was
363 lower.

364 3.3 Effect of dependent scattering on the bulk scattering coefficient in raw and homogenized 365 milk samples

366 To test the validity of the independent scattering condition, the BOP for the raw milk
367 sample and the milk samples homogenized for 120 and 1200 s were, next to their pure state
368 ($1/1$), also measured after a two-fold ($1/2$) and three-fold ($1/3$) dilution with deionized water. The
369 concentration of scattering particles in the pure, two-fold and three-fold diluted samples was
370 respectively 6.52%, 3.26% and 2.17% (v/v). No consistent differences between the g spectra
371 for the 3 particle concentrations could be detected for any of the 3 samples (results not
372 shown). This is probably because the noise on the measured g spectra was more significant
373 than the differences caused by the effect of dependent scattering at these particle volume
374 concentrations. It is well known that an anisotropy factor derived from DIS and unscattered
375 transmittance measurements is more prone to measurement noise in comparison to the other
376 derived BOPs [88]. Moreover, the effect of dependent scattering on the anisotropy factor is
377 only significant for relatively high volume concentrations of the scattering particles [48].

378 In Fig. 9, the measured μ_s spectra for the 3 samples (Raw, 120 s and 1200 s) and the 3
379 considered particle concentrations ($1/1$, $1/2$ and $1/3$) are shown. The μ_s spectra for the ‘ $1/2$ ’ and
380 ‘ $1/3$ ’ samples were multiplied with respectively 2 and 3 to scale them all to the same particle
381 concentration according to the linear independent scattering relation. Consequently, the
382 differences between the scaled μ_s spectra for the same sample are caused by dependent
383 scattering. For all samples, the highest volume concentration of scattering particles (pure
384 sample) clearly produced lower μ_s spectra as could be expected from the linear independent
385 scattering relation [Fig. 9]. Moreover, this effect was, to a smaller extent, still present in the μ_s
386 spectra for the two-fold dilution of the homogenized samples. For the raw milk sample, no
387 differences ($R^2 = 0.999$) were found between the scaled μ_s spectra for the ‘ $1/2$ ’ and ‘ $1/3$ ’
388 samples. This indicates that the independent scattering relation is valid up till higher particle
389 volume concentrations if the scattering particles are larger and confirms the observations
390 reported by other researchers [48,88,95–99]. Therefore, in raw milk samples, independent
391 scattering is valid up to particle volume concentrations of at least 3.26% (two-fold dilution).
392 For the homogenized samples, the difference between the scaled μ_s spectra for the ‘ $1/2$ ’ and
393 ‘ $1/3$ ’ samples is very small. Therefore, dependent scattering is likely absent in homogenized
394 milk samples with particle volume concentrations up to 2.17% (three-fold dilution).

395 3.4 Simulation of the bulk optical properties in raw and homogenized milk samples

396 As dependent scattering was present in the pure milk samples and in most of the two-fold
397 diluted samples [Fig. 9], while the simulation algorithm assumes independent scattering, the
398 BOP were simulated and validated for the three-fold diluted samples. The BOP were
399 simulated separately for the casein micelle fraction and fat globules, and were recombined
400 afterwards to obtain the BOP of the milk samples. The simulated bulk absorption coefficient
401 spectra matched well ($R^2 \geq 0.973$) with the measured ones. These spectra are, however, not
402 shown here as there was no visible difference between them and the μ_a spectrum of their
403 water fraction (96.6% v/v for three-fold diluted milk) [90]. The simulated and measured
404 spectra for respectively the bulk scattering coefficient spectra (μ_s) and scattering anisotropy
405 spectra (g) are shown in Fig. 10 and Fig. 11. As illustrated in these figures and Table 2, the
406 agreement between the simulated and measured bulk scattering properties for the three-fold
407 diluted samples is very good. The poorest agreement ($R^2 = 0.777$) was found between the
408 simulated and measured g spectra for the raw milk sample. As the anisotropy factor for raw
409 milk is nearly independent of the radiation wavelength, the very small disagreement (RMSD =
410 2.33×10^{-3}) already resulted in a rather poor R^2 . Comparison between all other simulated and
411 measured bulk scattering properties resulted in an R^2 above 0.94. For all three samples
412 together (raw, 120 s and 1200 s), the correlation between the simulated and measured

413 properties was very high, indicated by the very high coefficients of determination
414 (respectively 0.990 and 0.996) for the bulk scattering coefficients and the anisotropy factors
415 [Table 2]. The effect of homogenization on the measured and simulated bulk scattering
416 properties is fairly the same. This indicates the high potential of the Mie solution, after
417 inversion, to accurately estimate the PSD from BOP measurements. The high correlation
418 between simulated and measured bulk optical properties indicates a successful triple
419 validation of (1) the PSD measurements, (2) the BOP measurements and (3) the simulation
420 algorithm. Small differences between the simulated and measured values in Fig. 10 and Fig.
421 11 could be due to errors in the measurement of the BOP and/or PSDs and/or because of
422 invalidity of the perfect sphericity and independent scattering assumptions. However, they are
423 more likely the result of incorrect simulation-inputs for the refractive indices. Therefore,
424 accurate measurement of the complex refractive indices for milk fat, casein and milk serum
425 for the entire Vis/NIR wavelength range could enable more accurate simulations.

426 In contrast to the preliminary results presented in [40], the present manuscript shows a
427 much better agreement between simulated and measured BOP. Likely, this is the result of
428 optimized unscattered transmittance measurements and improved sensitivity for submicron
429 particles in the PSD measurements.

430 The simulated bulk scattering coefficient spectrum and scattering anisotropy spectrum for
431 the ‘virtual’ casein fraction are also shown in respectively Fig. 10 and Fig. 11. As the casein
432 micelles are relatively small [Table 1(b)], the g spectrum is low and is strongly dependent on
433 the radiation wavelength. Moreover, in combination with a low volume concentration of
434 casein, relative to the fat fraction (paragraph 2.4), this results in a very low μ_s spectrum, with
435 nearly no scattering in the NIR. As a consequence, the μ_s spectrum for the casein micelles will
436 not contribute much to the μ_s spectra of the milk samples. Likewise, also the simulated g
437 spectrum for the casein micelles will have no substantial impact on the g spectra of the raw
438 and homogenized milk samples.

439 Casein micelles are, however, not perfectly spherical, but a nearly-spherical cluster of
440 smaller casein sub-micelles [5]. For non-spherical particles, the simulated μ_a and μ_s for
441 equivalent spheres will be sufficiently accurate. However, the anisotropy factor g and
442 scattering phase function $p(\theta)$ can differ significantly [37]. For this reason, the simulated g
443 spectrum for the casein micelles in Fig. 11 might not be entirely correct. The T -matrix method
444 for sphere clusters could be a good alternative for the Mie solution [100–102]. However, as it
445 introduces extra degrees of freedom, it might import additional sources of error and should,
446 therefore, be validated thoroughly. As the fat globules dominate the bulk scattering properties
447 of milk, it is difficult to study the BOP of the casein micelles separately. Therefore, isolation
448 of the casein micelles without damaging and without loss of micelles is essential in order to
449 measure their exact BOP and validate the simulated results.

450 To enable online PSD measurements on turbid colloids with a particle volume
451 concentration above 2 – 4%, dilution is not preferred to reduce dependent scattering.
452 Therefore, a model is needed that describes the effect of dependent scattering on the bulk
453 scattering properties. Recently, the Percus-Yevick model for dependent scattering was
454 proposed and successfully validated on colloidal suspensions [49,50,103,104]. Moreover,
455 studies on milk showed that this model is accurate for total volume fraction of fat globules
456 and/or casein micelles up to 45% [49,50]. Further exploration of this model is, however,
457 needed to evaluate the influence of the radiation wavelength, particle and host refractive
458 indices and particle geometry, size, polydispersity and density [23,25].

459 Online measurement of the BOP for an undiluted colloidal suspension is possible with
460 time- or spatially-resolved Vis/NIR spectroscopy [38–42]. With the development of a general
461 model that describes dependent scattering, measured bulk scattering properties could be
462 transformed to follow the independent scattering relation. Consequently, after inversion of the
463 Mie solution, the transformed bulk scattering properties could be consulted to predict the PSD
464 of the suspended particles [105].

465 **4. Conclusions**

466 Ultrasonic homogenization of milk reduces the size of its fat globules in relation to the
467 exposure time. As a result, the bulk scattering properties change significantly and can,
468 therefore, be consulted to extract size distribution information. With decreasing diameter of
469 the fat globules, the visible (Vis) and near-infrared (NIR) bulk scattering coefficient and
470 scattering anisotropy factor reduce and become more dependent on the radiation wavelength.

471 In absence of dependent scattering, the Mie solution can accurately ($R^2 \geq 0.990$) describe
472 the relation between the fat globule size and the bulk scattering properties of a milk sample.
473 Therefore, inversion of the Mie solution could allow for the accurate estimation of the size
474 distribution from measured bulk scattering properties. As this would enable the prediction of
475 microstructure-related quality parameters, next to the compositional analysis used nowadays,
476 this would considerably increase the value of Vis/NIR spectroscopy for agriculture and food
477 industry. However, before this approach can be executed on very turbid suspensions like pure
478 milk, further research is required to develop a model that describes dependent scattering in
479 such media. Additionally, more research is needed with respect to the inversion of the Mie
480 scattering model in order to enable fast, accurate and robust prediction of the particle size
481 from bulk scattering properties.

482 Finally, it was found that casein micelles only have a limited contribution to the Vis/NIR
483 bulk scattering properties of milk. Nevertheless, better understanding of the scattering by
484 casein micelles would further improve the model that related the fat globule size with the
485 Vis/NIR bulk scattering properties of the milk sample.

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725

1 **Figures**

2

3 Fig. 1 The cumulative volume-frequency particle size distributions for the raw and
4 homogenized milk samples.

5

6 Fig. 2 The cumulative volume-frequency particle size distributions for the raw and
7 homogenized milk samples after subtraction of the casein contribution. The casein fraction
8 itself is also illustrated.

9

10 Fig. 3 Microscopic images (1250x magnification) of fat globules in raw milk. The black scale
11 bar in the bottom-right corner represents 10 μm .

12

13 Fig. 4 Microscopic images (1250x magnification) of fat globules in the raw milk sample (Fig.
14 3) after 120 s of ultrasonic homogenization. The black scale bar in the bottom-right corner
15 represents 10 μm .

16

17 Fig. 5 Microscopic images (1250x magnification) of fat globules in the raw milk sample (Fig.
18 3) after 1200 s of ultrasonic homogenization. The black scale bar in the bottom-right corner
19 represents 10 μm .

20

21 Fig. 6 Measured bulk absorption coefficient spectra (μ_a) for the undiluted raw and
22 homogenized milk samples.

23

24 Fig. 7 Measured bulk scattering coefficient spectra (μ_s) for the undiluted raw and
25 homogenized milk samples.

26

27 Fig. 8 Measured scattering anisotropy spectra (g) for the undiluted raw and homogenized milk
28 samples.

29

30 Fig. 9 Bulk scattering coefficient μ_s for raw and homogenized (for 120 s and 1200 s) milk
31 samples measured on pure samples ($1/1$) and two- ($1/2$) and three-fold dilutions ($1/3$) of the same
32 sample. The μ_s spectra for the diluted samples ($1/2$ and $1/3$) are rescaled to the particle
33 concentration of the pure sample through multiplication by respectively 2 and 3, following the
34 linear independent scattering relation.

35

36 Fig. 10 Measured and simulated (*) bulk scattering coefficient spectra (μ_s) for three-fold
37 diluted raw and homogenized milk samples. The simulated bulk scattering coefficient
38 spectrum for the 'virtual' casein fraction is also illustrated.

39

40 Fig. 11 Measured and simulated (*) scattering anisotropy spectra (g) for three-fold diluted raw
41 and homogenized milk samples. The simulated scattering anisotropy spectrum for the 'virtual'
42 casein fraction is also illustrated.

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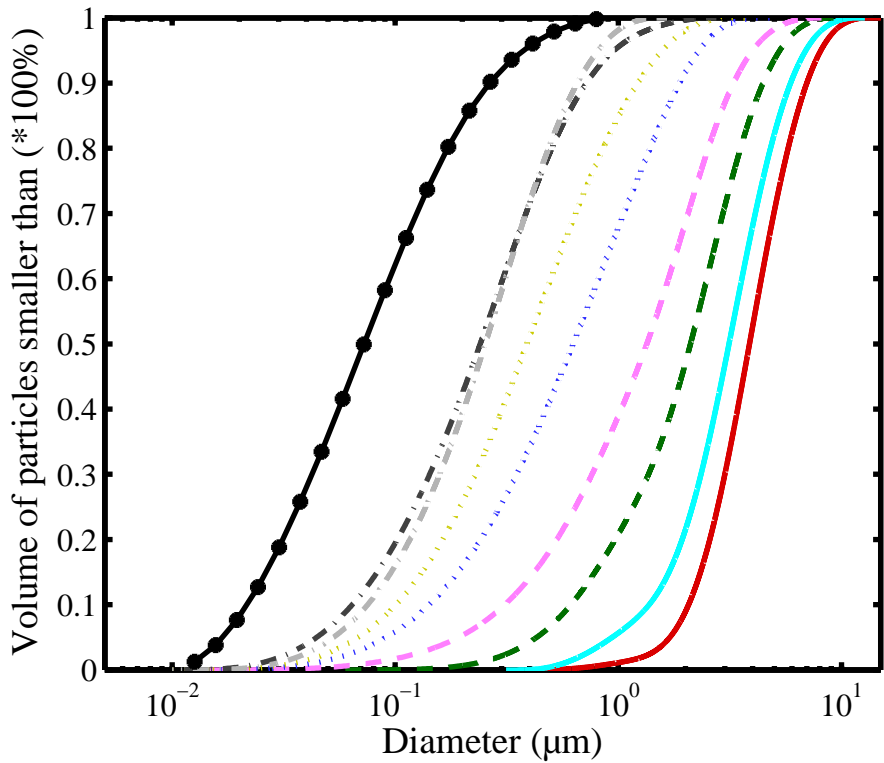
44 **Tables**

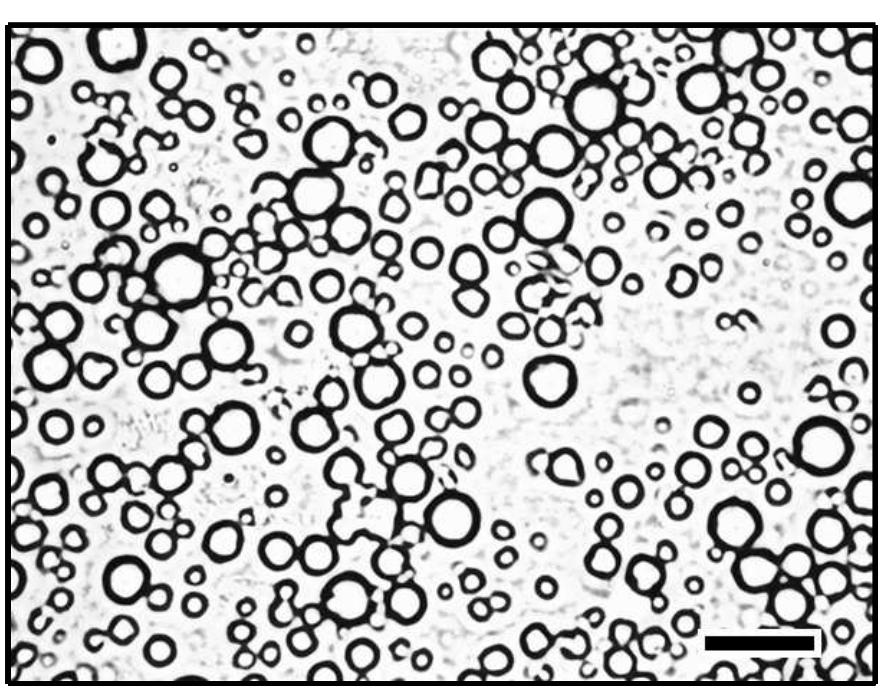
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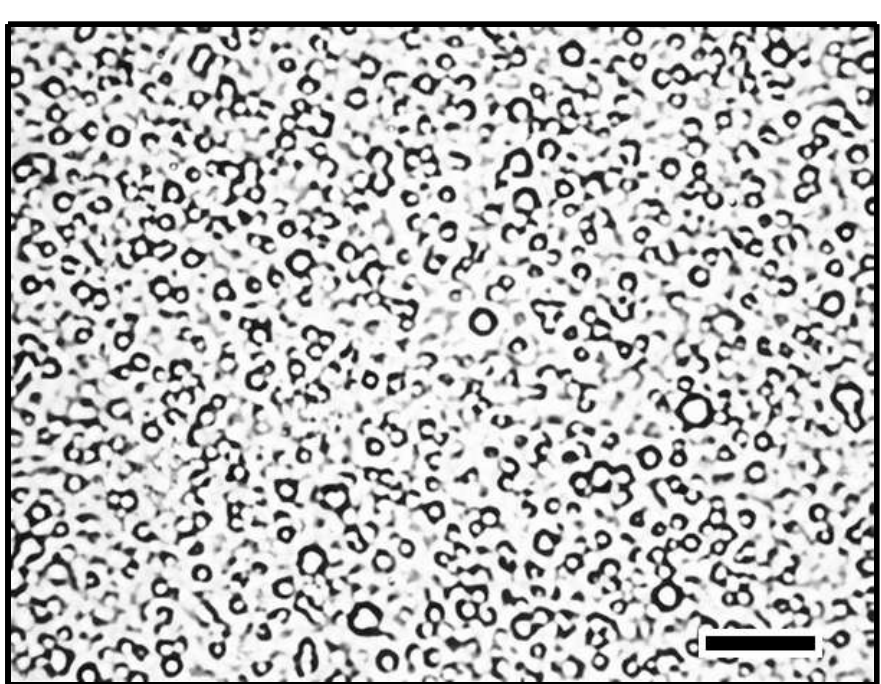
46 Table 1 Most important properties of the volume-frequency particle size distributions for the
47 raw (Raw), homogenized (* s) and skimmed (Skim) milk samples (Fig. 1 and Fig. 2) before
48 (a) and after (b) subtraction of the casein micelle contribution. The properties of the casein
49 micelle fraction (Casein) itself are also shown in (b). The Dv10, Dv50 and Dv90 are
50 respectively the 10th, 50th and 90th percentile of the volume-based particle size distribution.
51 The D32 and D43 are respectively the Sauter and DeBroukere mean particle diameter.

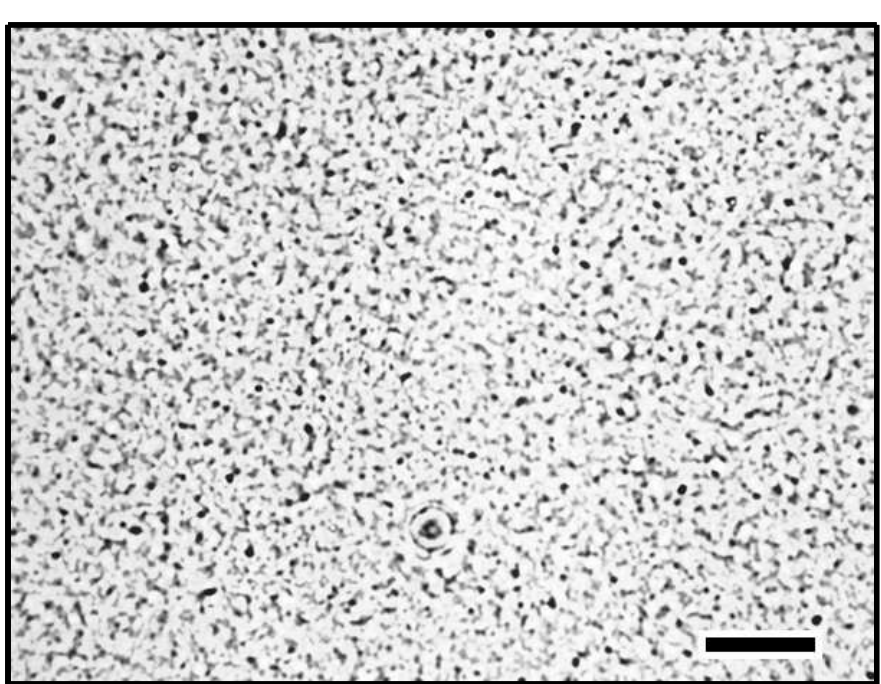
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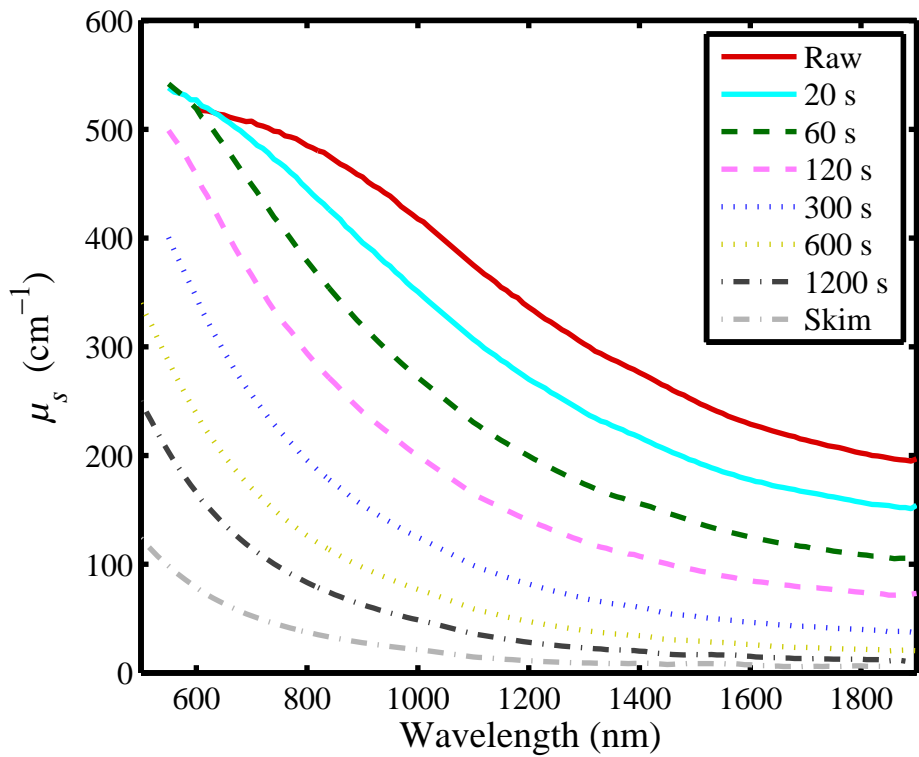
53 Table 2 Coefficients of determination, indicating the agreement between measured and
54 simulated bulk optical properties for the raw and homogenized milk samples illustrated in Fig.
55 10 and Fig. 11.

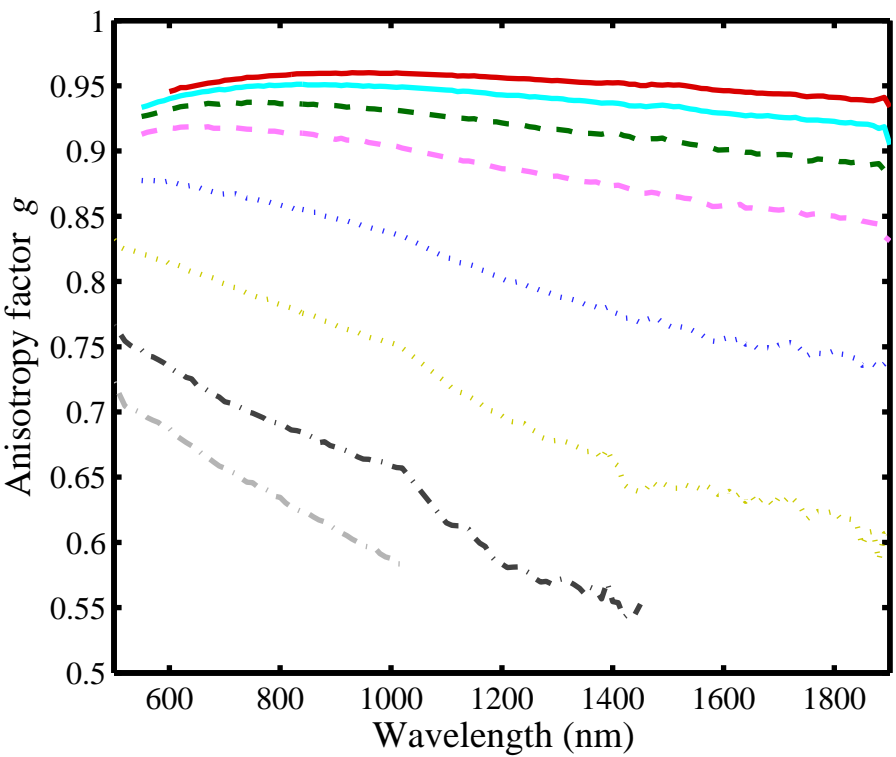


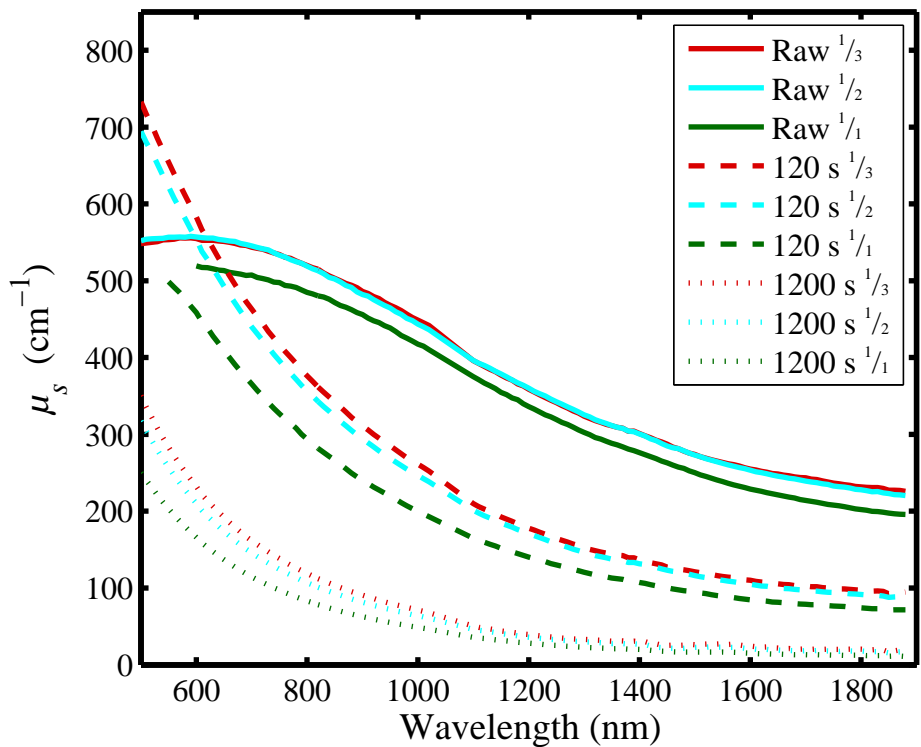


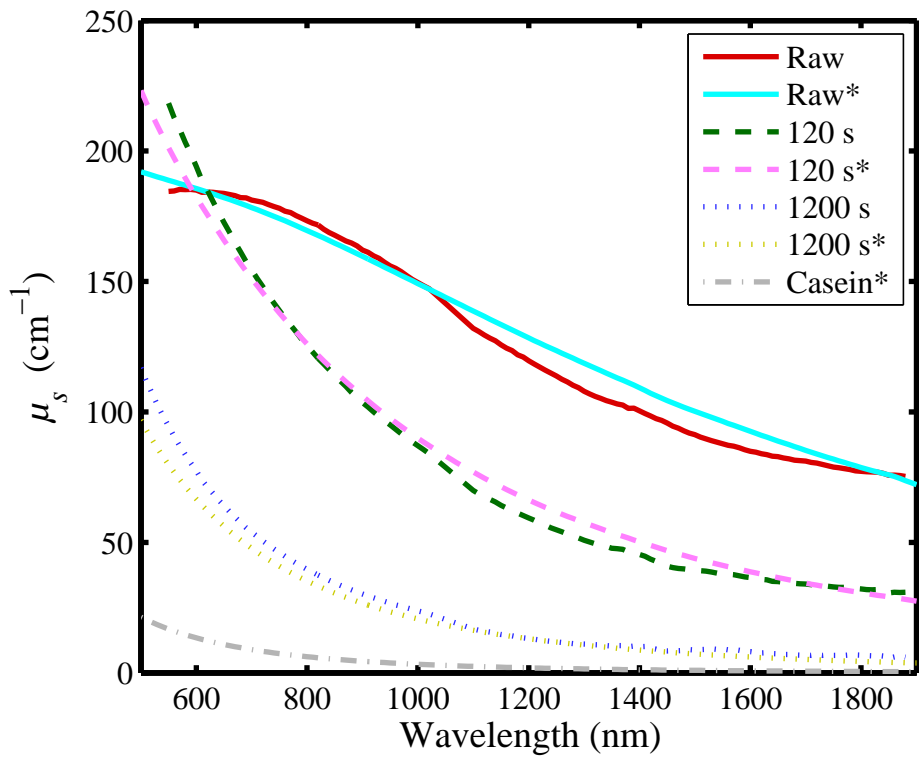












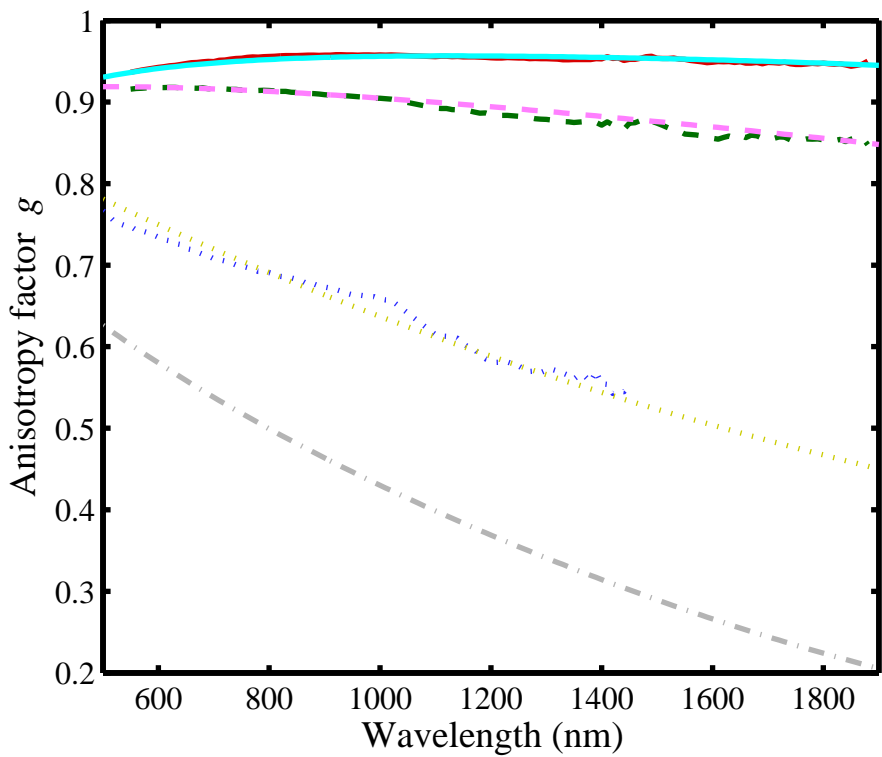


Table 1 Most important properties of the volume-frequency particle size distributions for the raw (Raw), homogenized (* s) and skimmed (Skim) milk samples (Fig. 1 and Fig. 2) before (a) and after (b) subtraction of the casein micelle contribution. The properties of the casein micelle fraction (Casein) itself are also shown in (b). The Dv10, Dv50 and Dv90 are respectively the 10th, 50th and 90th percentile of the volume-based particle size distribution. The D32 and D43 are respectively the Sauter and DeBroukere mean particle diameter.

Sample	(a) With casein (μm)					(b) Casein subtraction (μm)				
	Dv10	Dv50	Dv90	D32	D43	Dv10	Dv50	Dv90	D32	D43
Raw	0.0306	0.411	5.69	0.0947	2.09	2.06	3.87	6.99	3.34	4.24
20 s	0.0295	0.312	4.67	0.0897	1.65	1.37	3.13	5.91	2.44	3.44
60 s	0.0288	0.222	3.30	0.0848	1.07	0.613	2.11	4.45	1.31	2.36
120 s	0.0288	0.195	2.35	0.0826	0.753	0.310	1.34	3.28	0.648	1.61
300 s	0.0288	0.165	1.27	0.0789	0.431	0.142	0.627	1.88	0.321	0.844
600 s	0.0285	0.143	0.777	0.0748	0.298	0.101	0.389	1.22	0.226	0.545
1200 s	0.0287	0.127	0.518	0.0721	0.219	0.0642	0.240	0.719	0.145	0.336
Skim	0.0248	0.0984	0.407	0.0606	0.167	0.0753	0.253	0.655	0.162	0.317
Casein						0.0217	0.0724	0.264	0.0501	0.116

Table 2 Coefficients of determination, indicating the agreement between measured and simulated bulk optical properties for the raw and homogenized milk samples illustrated in Fig. 10 and Fig. 11.

Sample	$R_{\mu_s}^2$	R_g^2
Raw	0.979	0.777
120 s	0.990	0.943
1200 s	0.954	0.964
Combined	0.990	0.996